Specific heat (definition) and its relation to the temperature derivative of total internal energy

Change of the heat Q consists of the change of total energy + the work done:

$$\frac{\mathrm{d}Q}{\mathrm{d}t} = \frac{\mathrm{d}E}{\mathrm{d}t} + P\frac{\mathrm{d}V}{\mathrm{d}t} , \quad \frac{\mathrm{d}E}{\mathrm{d}t} = T\frac{\mathrm{d}S}{\mathrm{d}t} - P\frac{\mathrm{d}V}{\mathrm{d}t} , \quad \mathrm{d}Q/\mathrm{d}t = T \,\mathrm{d}S/\mathrm{d}t.$$

Hence, the specific $C_v = T(\partial S/\partial T)_V$, heat C=dQ/dT is $C_p = T(\partial S/\partial T)_P$.

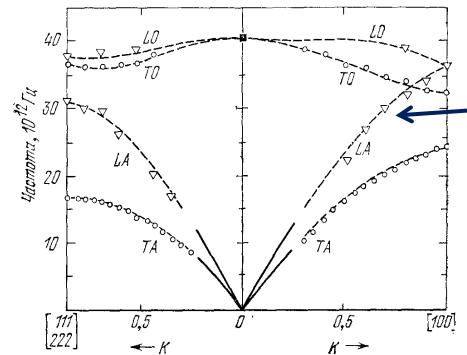
The energy differential $dE = T dS - P dV \implies C_v = (\partial E/\partial T)_V$

Phonon contribution to the total energy

The energy of all excitations (phonons) is given by the sum over all quantum states, which are numbered by the wave vector $k=p/\hbar$:

$$E_{ph}(T) = \sum_{k,\alpha} \varepsilon_{\alpha}(k) n_{k}(\varepsilon) = \sum_{\alpha} \int \frac{Vd^{3}k}{(2\pi)^{3}} \varepsilon_{\alpha}(k) n_{k}(\varepsilon_{\alpha})$$

The filling number n_k of the quantum states of phonons is given by Bose-Einstein distribution function:



$$n_k(\varepsilon) = \frac{1}{\exp([\varepsilon(k) - \mu]/k_B T) - 1}$$

The phonon dispersion $\omega(k)$ may consist of several branches α .

The phase volume $\int Vd^3k/(2\pi)^3$ gives the number of quantum states.

Specific heat (heat capacity of unit mass)
$$C_v = (\partial E/\partial T)_V$$

Heat capacity and density of phonon sates in 1D

Energy of phonon gas
$$E = \sum_{K} \left\langle n_{K} \right\rangle \hbar \omega_{K}$$
, or $E = \int d\omega \, \mathcal{D} \left(\omega \right) \left\langle n \left(\omega, T \right) \right\rangle \hbar \omega$.

In 1D case $\mathcal{D} \left(\omega \right) \, d\omega = \frac{L}{\pi} \frac{dK}{d\omega} \, d\omega = \frac{L}{\pi} \frac{d\omega}{d\omega/dK}$. $\left\langle n \right\rangle = \frac{1}{\exp(\hbar \omega/\tau) - 1}$

- 1. In Debye approximation $\omega(K) = vK$, and $\mathscr{D}(\omega) = \frac{L}{\pi v}$ при $\omega \leqslant \frac{v\pi}{\sigma}$
- 2. In the Einstein model $\mathscr{D}(\omega) = N \delta(\omega \omega_E)$,
- 3. For 1D chain of atoms $\omega = \omega_{\text{max}} \left| \sin \frac{1}{2} Ka \right|$, $K = \frac{2}{a} \arcsin \frac{\omega}{\omega_{\text{max}}}$,

$$\frac{dK}{d\omega} = \frac{2}{a} \frac{1}{(\omega_{\max}^2 - \omega^2)^{1/2}} \quad \text{and the density of phonon states} \quad \mathcal{D}(\omega) = \frac{L}{\pi} \frac{dK}{d\omega} = \frac{2L}{\pi a} \frac{1}{(\omega_{\max}^2 - \omega^2)^{1/2}}$$

$$\frac{d}{d\omega} = \frac{2}{a} \frac{1}{(\omega_{\max}^2 - \omega^2)^{1/2}} \quad \text{of phonon states} \quad \mathcal{D}(\omega) = \frac{L}{\pi} \frac{dK}{d\omega} = \frac{2L}{\pi a} \frac{1}{(\omega_{\max}^2 - \omega^2)^{1/2}}$$

$$\frac{d}{d\omega} = \frac{1}{\pi a} \frac{1}{(\omega_{\max}^2 - \omega^2)^{1/2}} \quad \text{density of states} \quad \text{in 1D}$$

$$\frac{dK}{d\omega} = \frac{2L}{\pi a} \frac{1}{(\omega_{\max}^2 - \omega^2)^{1/2}} \quad \text{density of states} \quad \text{in 1D}$$

$$\frac{dK}{d\omega} = \frac{2L}{\pi a} \frac{1}{(\omega_{\max}^2 - \omega^2)^{1/2}} \quad \text{density of states} \quad \text{density of states} \quad \text{in 1D}$$

$$\frac{dK}{d\omega} = \frac{2L}{\pi a} \frac{dK}{d\omega} = \frac{2L}{\pi a} \frac{1}{(\omega_{\max}^2 - \omega^2)^{1/2}} \quad \text{density of states} \quad \text{den$$

Density of phonon states in 3D (general case)

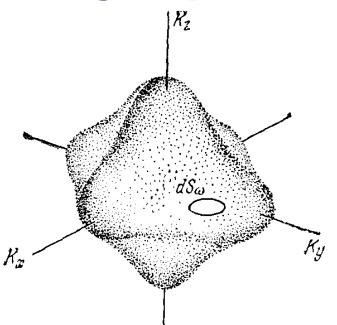


Рис. 6.10а. Элементарная площадка dS_{ω} на поверхности постоянной частоты в K-пространстве. Объем слоя между двумя поверхностями постоянной частоты ω и $\omega + d\omega$ равен $\int dS_{\omega} \, d\omega/(\mathrm{grad}_K \omega)$.

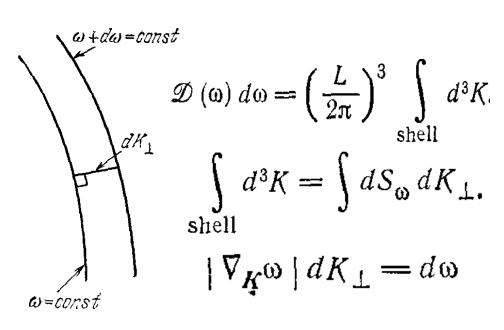
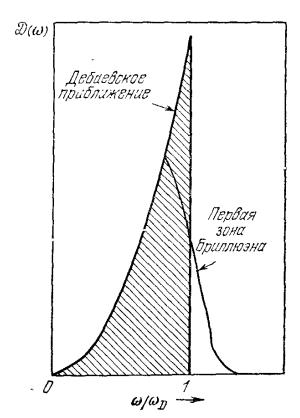


Рис. 6.10б. Величина dK_{\perp} есть расстояние между поверхностями постоянных частот ω и $\omega + d\omega$, взятое вдоль нормали к ним.

$$dS_{\omega} dK_{\perp} = dS_{\omega} \frac{d\omega}{|\nabla_{K}\omega|} = dS_{\omega} \frac{d\omega}{v_{g}}$$
, где $v_{g} = |\nabla_{K}\omega|$ — величина групповой скорости фонона.

$$\mathscr{D}(\omega) \ d\omega = \left(\frac{L}{2\pi}\right)^3 \int \frac{dS_{\omega}}{v_g} \ d\omega. \quad \Longrightarrow \quad \mathscr{D}(\omega) = \frac{V}{(2\pi)^3} \int \frac{dS_{\omega}}{v_g} \propto \omega^2$$

Density of phonon states in 3D Debye model



In Debye model [linear ω (k)] the number of phonon states with energy less than ω in a volume V=L³ is

$$N = \left(\frac{L}{2\pi}\right)^3 \frac{4\pi}{3} K^3 = \left(\frac{L}{2\pi}\right)^3 \frac{4\pi\omega^3}{3v^3} = \frac{V\omega^3}{6\pi^2 v^3}.$$

The Debye (cut off) frequency $\omega_D^3 = \frac{6\pi^2 v^3 N}{V}$

The thermal energy is given by

$$U = \int d\omega D(\omega) \langle n(\omega) \rangle \hbar \omega = \int_0^{\omega_D} d\omega \left(\frac{V\omega^2}{2\pi^2 v^3} \right) \left(\frac{\hbar \omega}{e^{\hbar \omega/\tau} - 1} \right)$$

$$U = \frac{3V\hbar}{2\pi^2 v^3} \int_0^{\omega_D} d\omega \frac{\omega^3}{e^{\hbar \omega/\tau} - 1} = \frac{3Vk_B^4 T^4}{2\pi^2 v^3 \hbar^3} \int_0^{x_D} dx \frac{x^3}{e^x - 1}$$

$$C_V = \frac{3V\hbar^2}{2\pi^2 v^3 k_B T^2} \int_0^{\omega_D} d\omega \frac{\omega^4 e^{\hbar \omega/\tau}}{(e^{\hbar \omega/\tau} - 1)^2} = 9Nk_B \left(\frac{T}{\theta} \right)^3 \int_0^{x_D} dx \frac{x^4 e^x}{(e^x - 1)^2}$$