



Infrared Hall Effect in cuprates*

H.D. Drew

Center for Nanophysics and Advanced Materials,
University of Maryland, College Park, MD

* Supported by NSF



Collaborators



Maryland

- D. Schmadel
 - L. Shi
 - G. Jenkins
 - R. Greene
-
- J. Černe, Buffalo
 - M. Houseknecht, Buffalo
-
- L. Rigal, Toulouse
 - A. Zimmers, Paris

Theory

- P. Coleman, Rutgers
- H. Kontani, Nagoya
- A. Millis, Columbia
- V. Yakovenko, Maryland

Materials

- I. Tsukada, CRIEPI, Tokyo
- Y. Ando, CRIEPI, Tokyo
- G. Gu, Brookhaven
- R. Hughes, McMaster
- J. Preston, McMaster



Outline



1. Introduction
2. IR Hall Effect
3. Cuprate phases
4. Optimally doped cuprates
5. Underdoped cuprates
6. Conclusions



IR Hall Effect: Hall spectroscopy

Magneto-Transport

Boltzmann theory

$$\sigma_{xx} = \frac{e^2}{hd} \oint dk \vec{\ell}_k$$

$$\sigma_{xy} = \frac{e^2}{hd} \frac{eB}{\hbar c} \oint dk \cdot \vec{\ell}_k \times d\vec{\ell}_k / dk$$

$$DC \quad \vec{\ell}_k = v_k \tau_k$$

$$IR \quad \vec{\ell}_k = \frac{v_k}{1/\tau_k - i\omega}$$

$$DC: \langle \ell_k \rangle$$

$$IR: \omega_H, \left\langle \frac{1}{\tau_k} \right\rangle$$

Sum Rules

IR Hall effect spectroscopy

$$\int_0^{\Omega_c} \frac{2}{\pi} \text{Re}(\sigma_{xx}) d\omega \equiv \omega_p^2 / 4\pi = \frac{e^2}{2} \sum_k \text{Tr}(m_k^{-1}) n_k$$

$$\int_0^{\Omega_c} \text{Re}(\sigma_{xy}) d\omega = 0$$

$$\int_0^{\Omega_c} \frac{2}{\pi} \text{Im}(\sigma_{xy}) \omega d\omega \equiv \omega_p^2 \omega_H / 4\pi = \frac{e^3 B}{2c_0} \sum_k \det(m_k^{-1}) n_k$$

$$\theta_H = \frac{\sigma_{xy}}{\sigma_{xx}}$$

$$\int_0^{\Omega_c} \frac{2}{\pi} \text{Re}(\theta_H) d\omega \equiv \omega_H = \frac{2eB}{c} \frac{\sum_k \det(m_k^{-1}) n_k}{\sum_k \text{Tr}(m_k^{-1}) n_k}$$

IR Hall effect:

The extended Drude model:

$$\sigma_{xx} = \frac{\omega_p^2}{4\pi} \frac{1}{\gamma - i\omega(1 + \lambda)}, \quad \sigma_{xy} = \frac{\omega_p^2}{4\pi} \frac{\omega_H}{(\gamma - i\omega(1 + \lambda))^2}$$

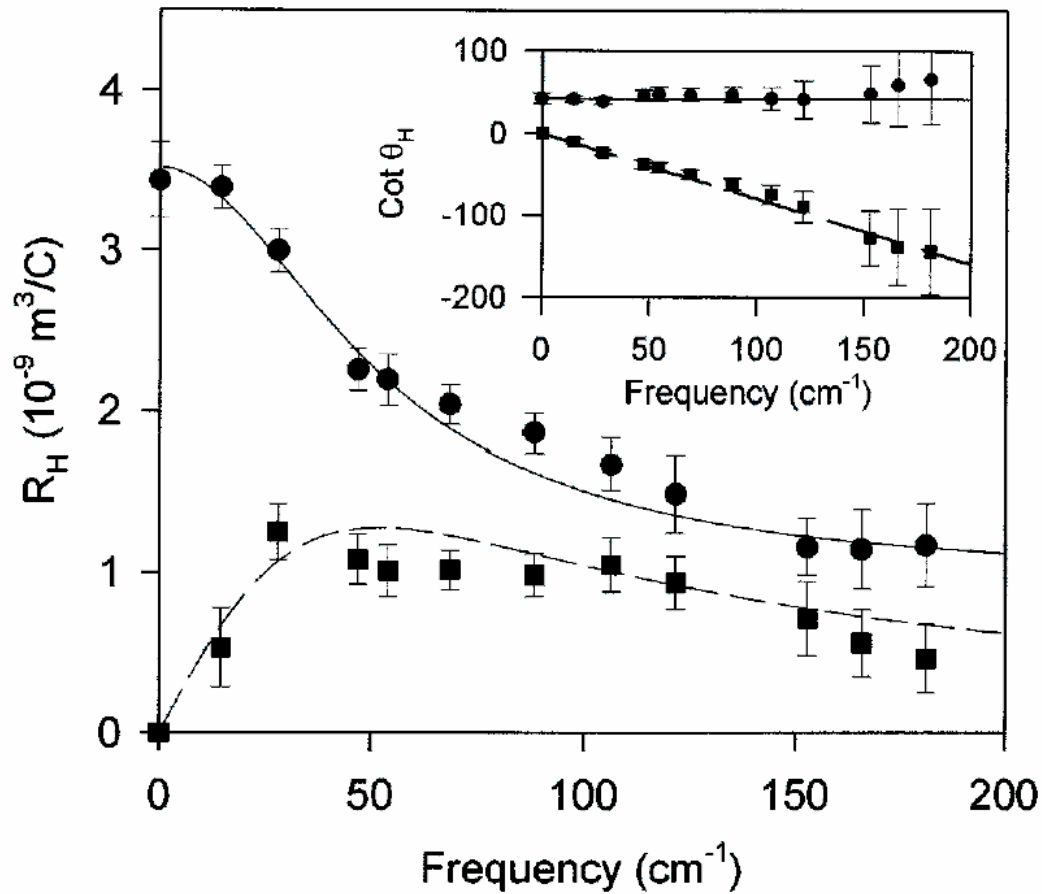
$$\theta_{xy} \equiv \frac{\sigma_{xy}}{\sigma_{xx}} = \frac{\omega_H}{\gamma - i\omega(1 + \lambda)}$$

$$R_H = \frac{1}{nec}$$

Optical self energy

$$\Sigma(\omega) \equiv \omega\lambda(\omega) + i\gamma(\omega)$$

R_H vs. ω YBCO₇



$$\sigma_{xx} = \frac{ne^2 / m}{\gamma - i\omega}$$

$$\sigma_{xy} = \frac{(ne^2 / m)(eB / mc)}{(\gamma - i\omega)^2}$$

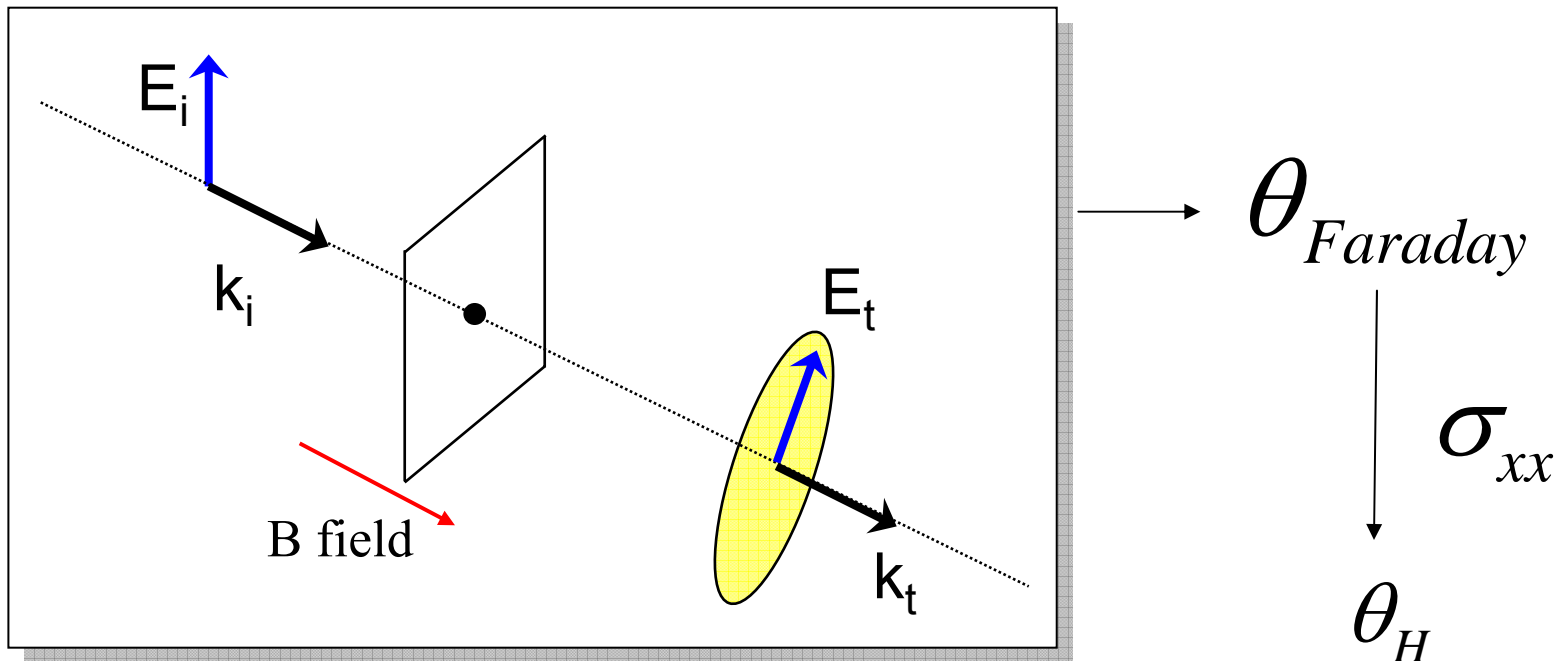
$$\rho_H = R_H B = \frac{\sigma_{xy}}{\sigma_{xx}^2} = \frac{B}{nec}$$

$$\theta_H = \frac{\sigma_{xy}}{\sigma_{xx}} = \frac{eB / mc}{\gamma - i\omega}$$

$$\gamma_{xy} \neq \gamma_{xx}$$

Kaplan *et al.*, PRL '96

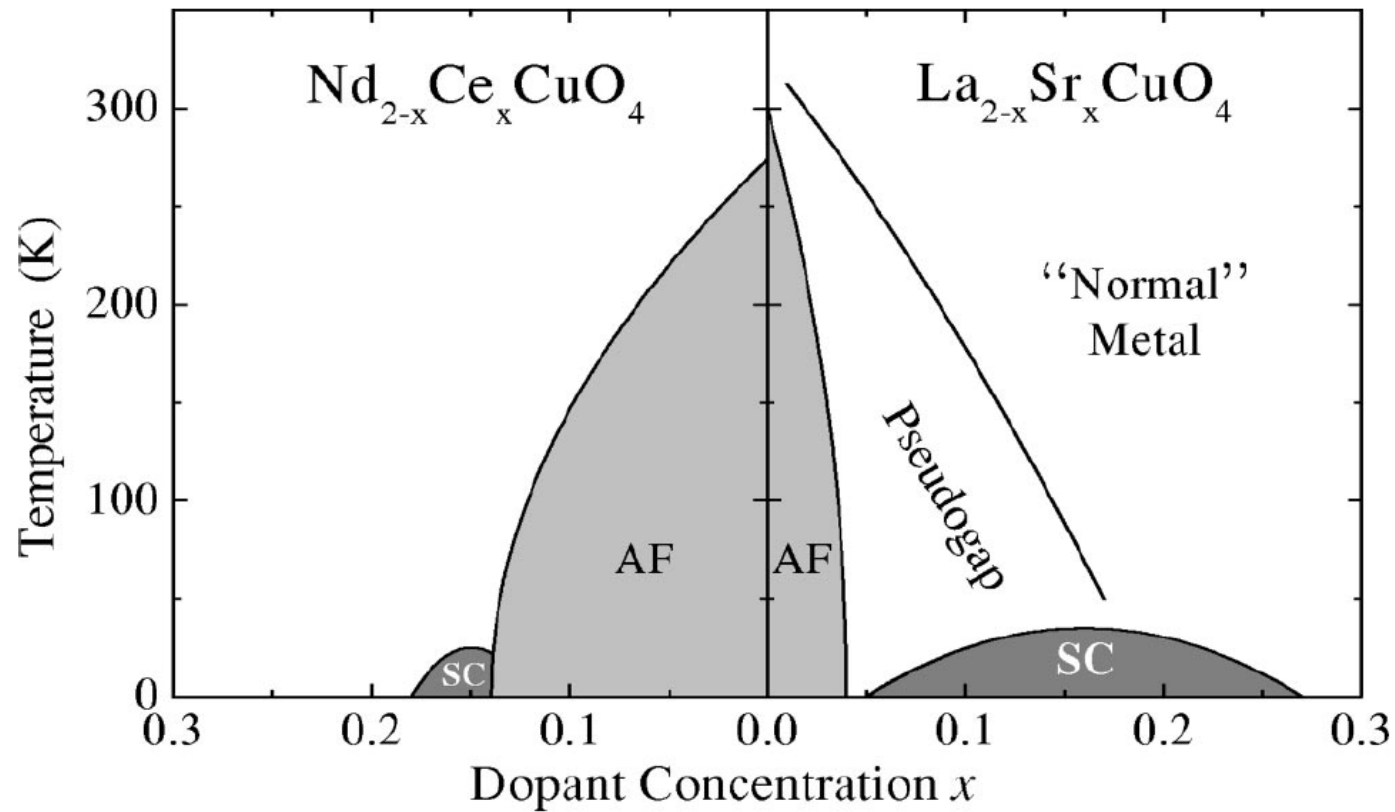
IR Hall Effect



$$\theta_H = \frac{\sigma_{xy}}{\sigma_{xx}}$$

J. Černe et al., RSI (2003)

Cuprate Phase Diagram



Fermi surfaces

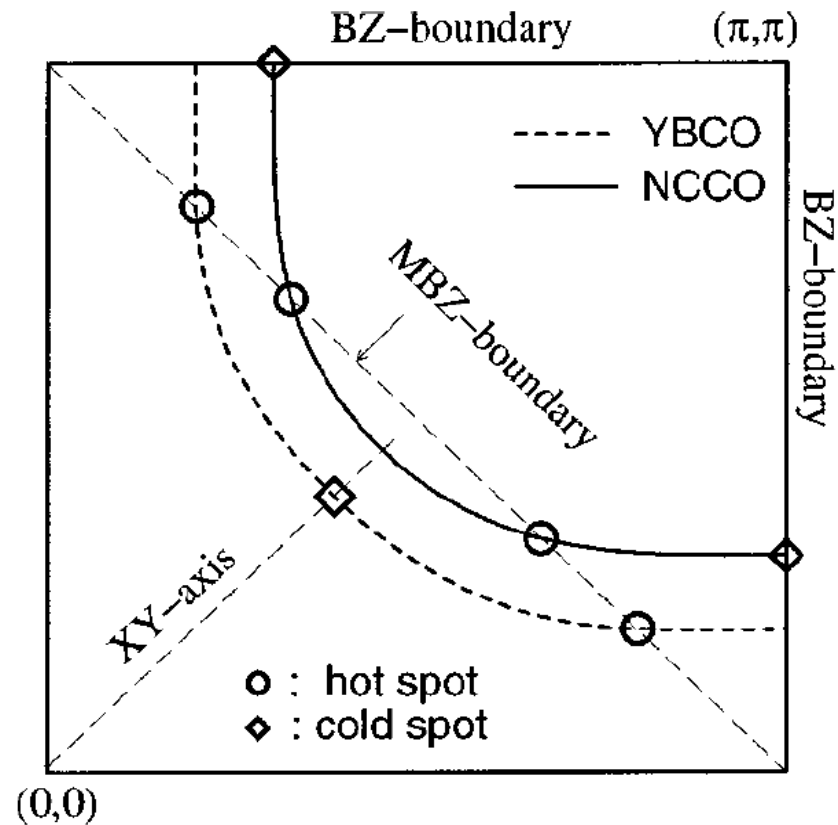


FIG. 4. The hot spots and the cold spots in YBCO and NCCO, respectively.

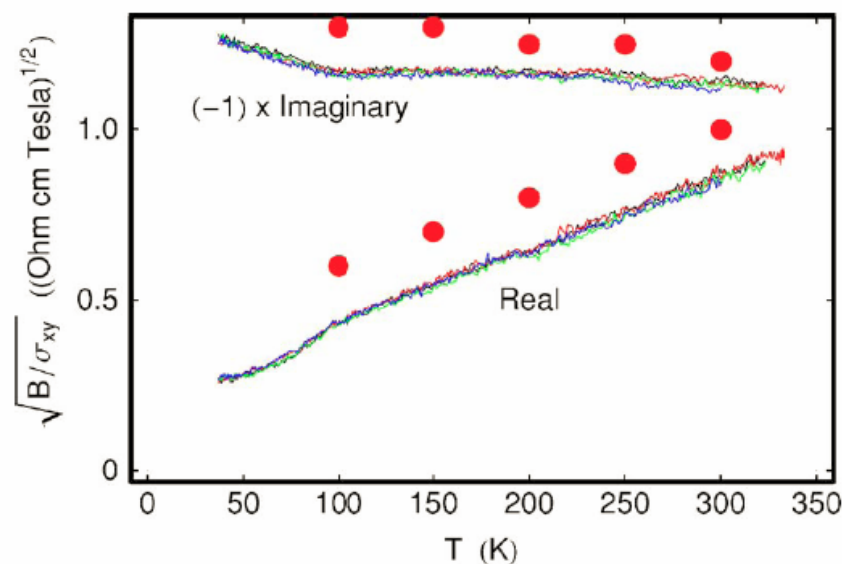
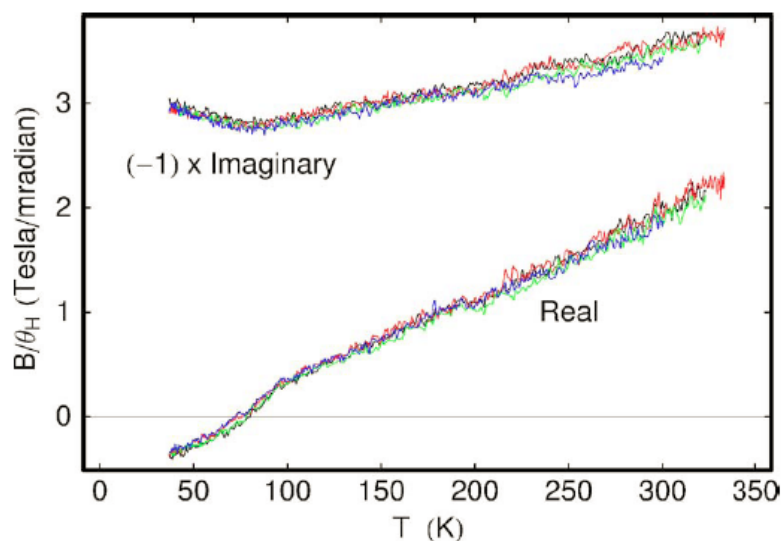
H. KONTANI, K. KANKI, AND K. UEDA PRB 59 (1999).

IR Hall Effect
in
Optimally doped Cuprates

Mid IR Hall in opt. Bi2212

Hall angle and Hall conductivity

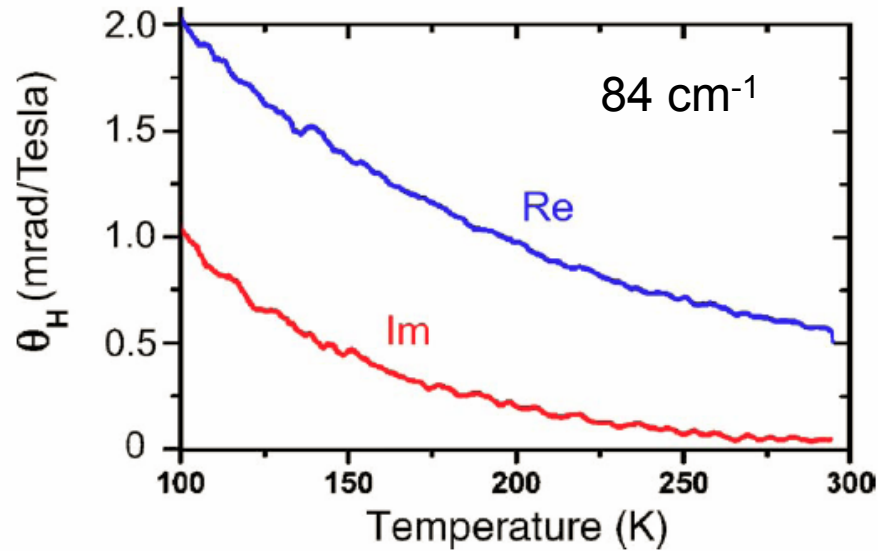
$$\omega = 1000 \text{ cm}^{-1}$$



$$\theta_H \equiv \frac{\sigma_{xy}}{\sigma_{xx}} = \frac{\omega_H}{\gamma - i\omega}, \quad \omega_H \equiv \frac{eB}{m_H c} \quad \sigma_{xy} = \frac{\omega_p^2}{4\pi} \frac{\omega_H}{(\gamma - i\omega)^2}$$

Schmadel, et al, PRB (2007)

Opt. Bi(2212): Far IR Hall



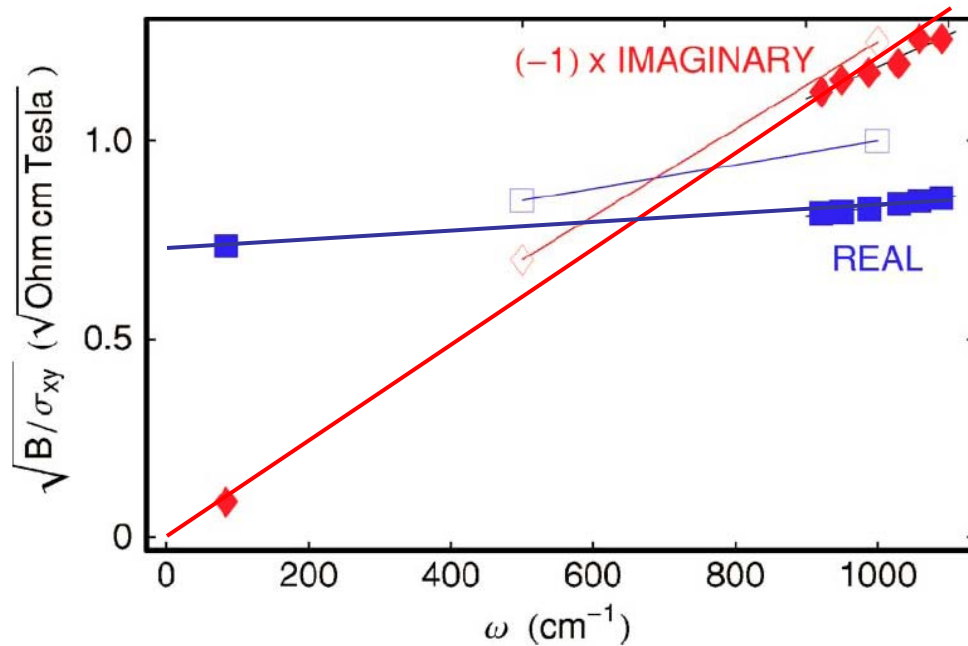
$$\theta_H = \frac{\omega_H}{\gamma - i\omega}$$

$$\omega_H \sim 0.3 \text{ cm}^{-1}/\text{T}$$

$$\gamma_H \sim aT^{1.65}$$

FIG. 3. (Color online) θ_H for optimally doped 2212 BSCCO versus temperature measured at 84 cm^{-1} averaged from five thermal scans and normalized to 1 T.

Optimally doped Bi2212



Drude model

$$\sigma_{xy}^{-1/2} = \frac{2}{\omega_p} \sqrt{\frac{\pi}{\omega_H}} (\gamma - i\omega(1 + \lambda))$$

$$\omega_H \sim 0.3 \text{ cm}^{-1}/T$$

Nearly Drude: inelastic scattering is weak in σ_{xy}

Schmadel et al., PRB (2007)

Hall and longitudinal scattering

Mid IR

	Bi2212
γ_{xx}^* , cm^{-1}	750
γ_{xy}^* , cm^{-1}	360
m_{xx}^*	2.0

$$\gamma^* = \frac{\gamma}{1 + \lambda} = \frac{\gamma}{m^* / m_b}$$

$$\omega = 1000 \text{ cm}^{-1}, T = 100 \text{ K}$$

Schmadel et al., PRB (2007)

Kontani theory vertex corrections to σ

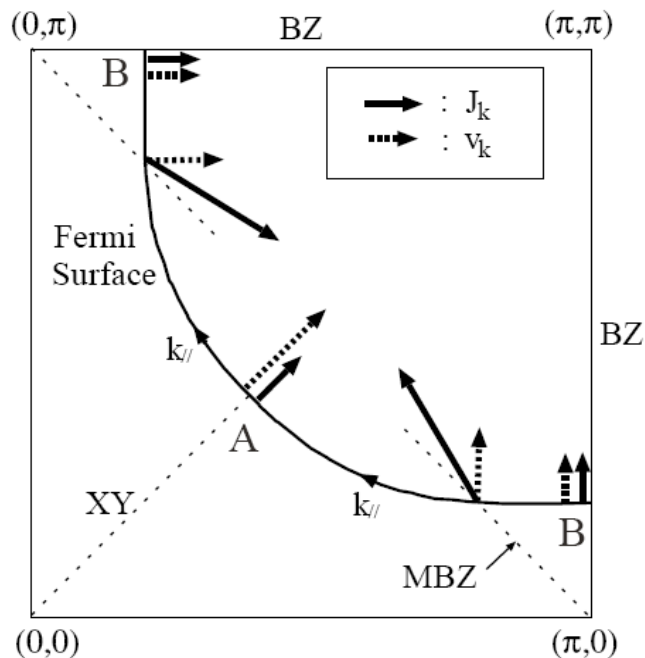
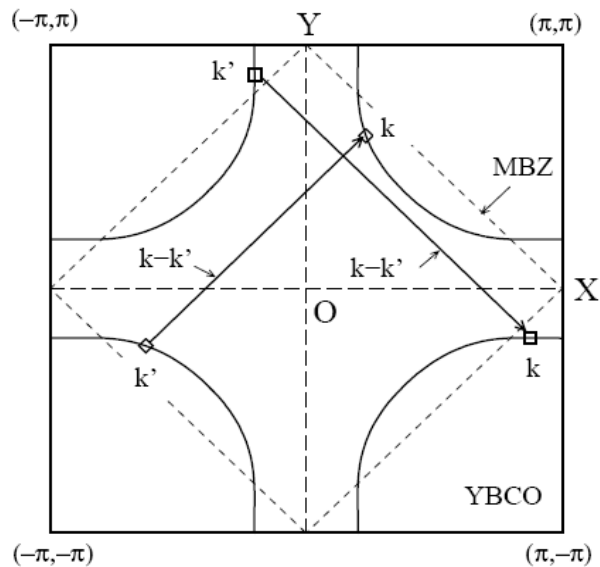
Antiferromagnetic fluctuations \rightarrow
FLEX approximation:

$$\chi_{\mathbf{k}}^s(\omega + i\delta) \approx \frac{\chi_{\mathbf{Q}}}{1 + \xi^2(\mathbf{q} - \mathbf{Q})^2 - i\omega/\omega_{\text{sf}}},$$

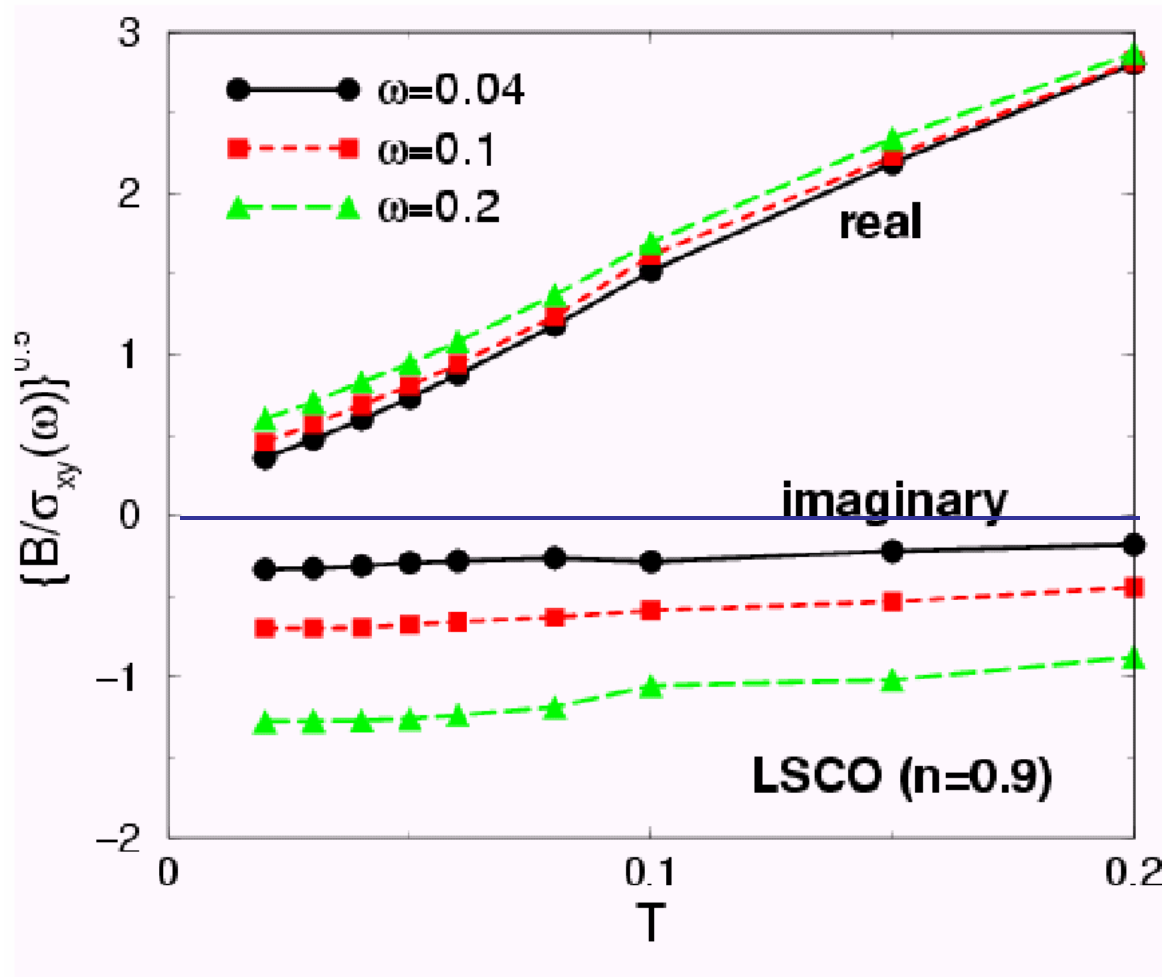
where ξ is the AF correlation length, and \mathbf{Q} is the nesting vector. $\chi_{\mathbf{Q}} \sim 1/\omega_{\text{sf}} \sim \xi^2$, and $\xi^2 \sim T^{-1}$ in the FLEX approximation,

**Conductivity calculated including
Current vertex corrections**

Kontani, J. Phys. Soc. Jpn. (2006)

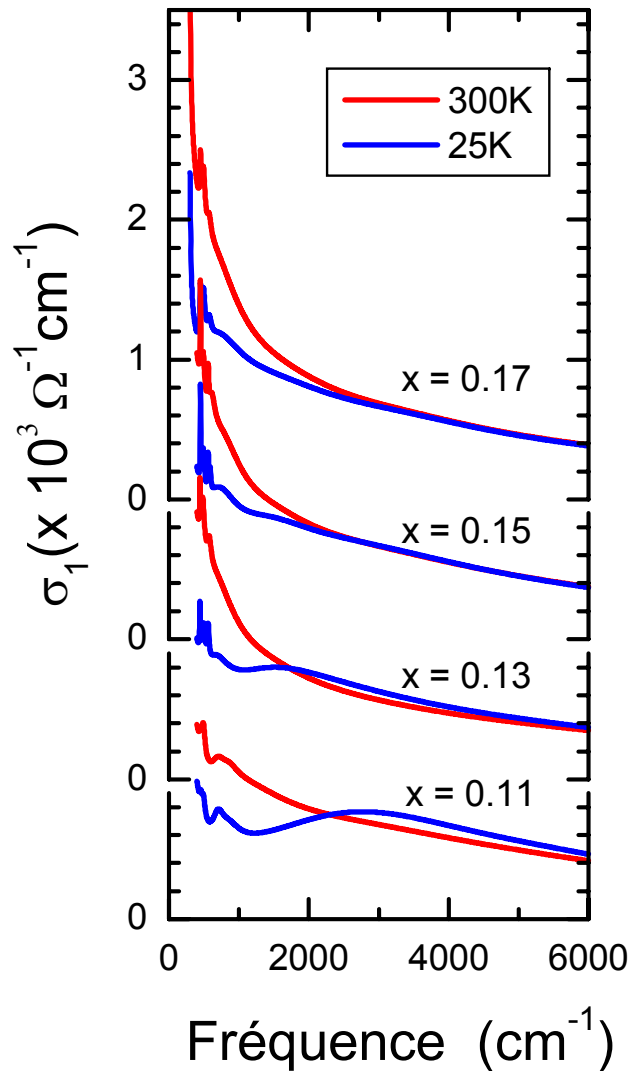


Kontani calculation: σ_{xy}



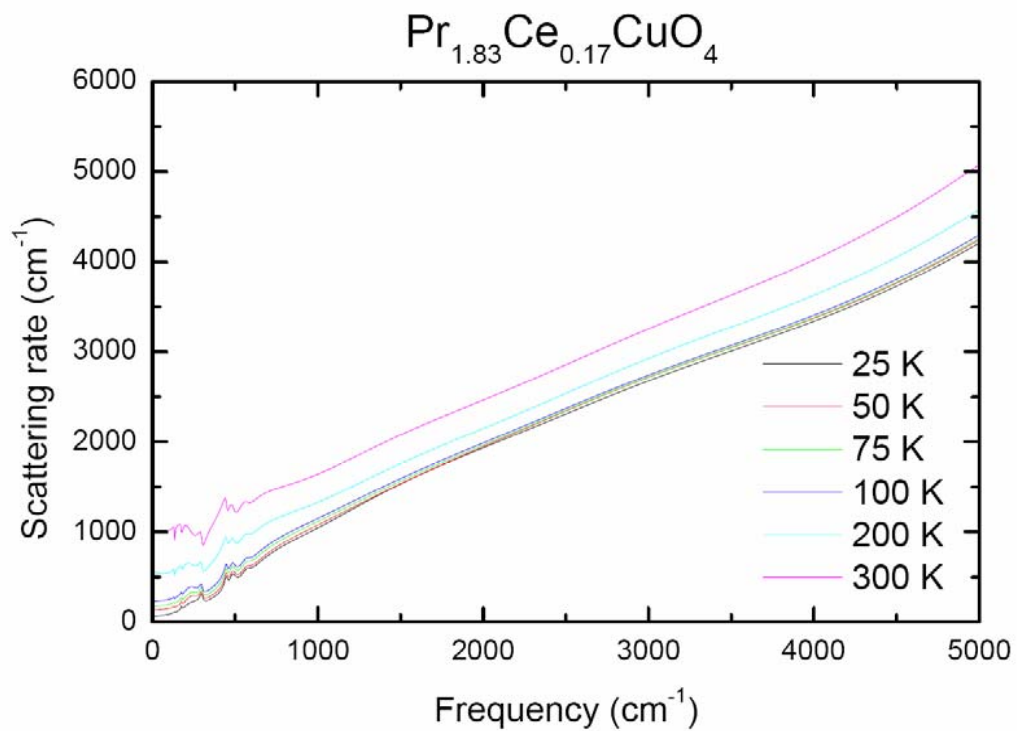
$$\sqrt{\sigma_{xy}^{-1}} = \frac{\gamma - i\omega(1 + \lambda)}{\sqrt{S_{xy}}}$$

e-doped: Optics of $\text{Pr}_{2-x}\text{Ce}_x\text{CO}_4$



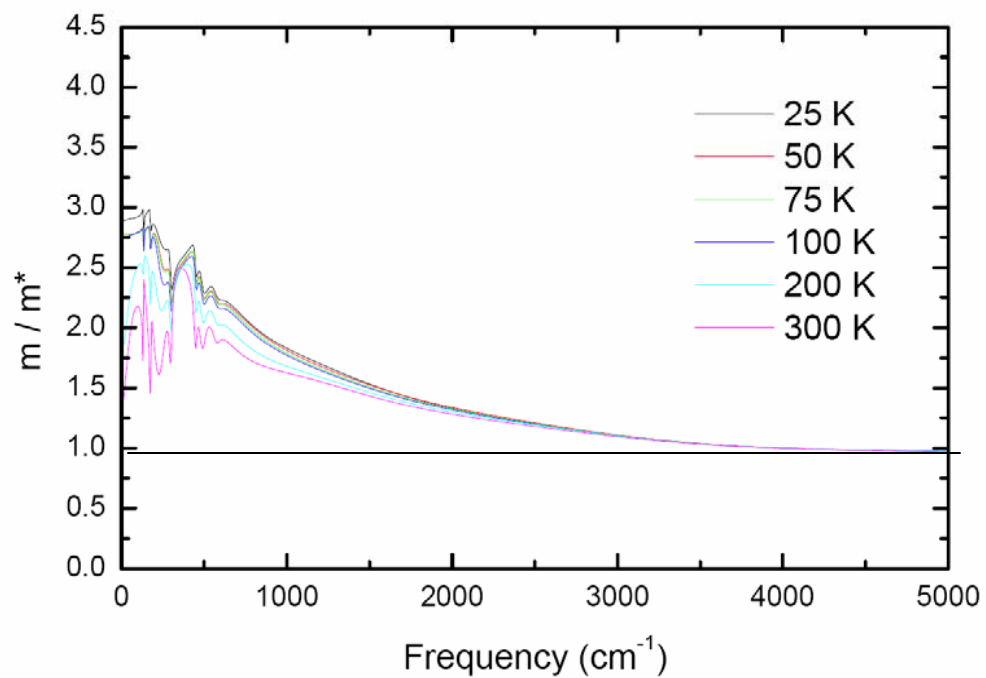
X=0.17
Extended
Drude analysis

$$\sigma_{xx}(\omega) = \frac{\omega_p^2 / 4\pi}{\gamma - i\omega(1 + \lambda)}$$



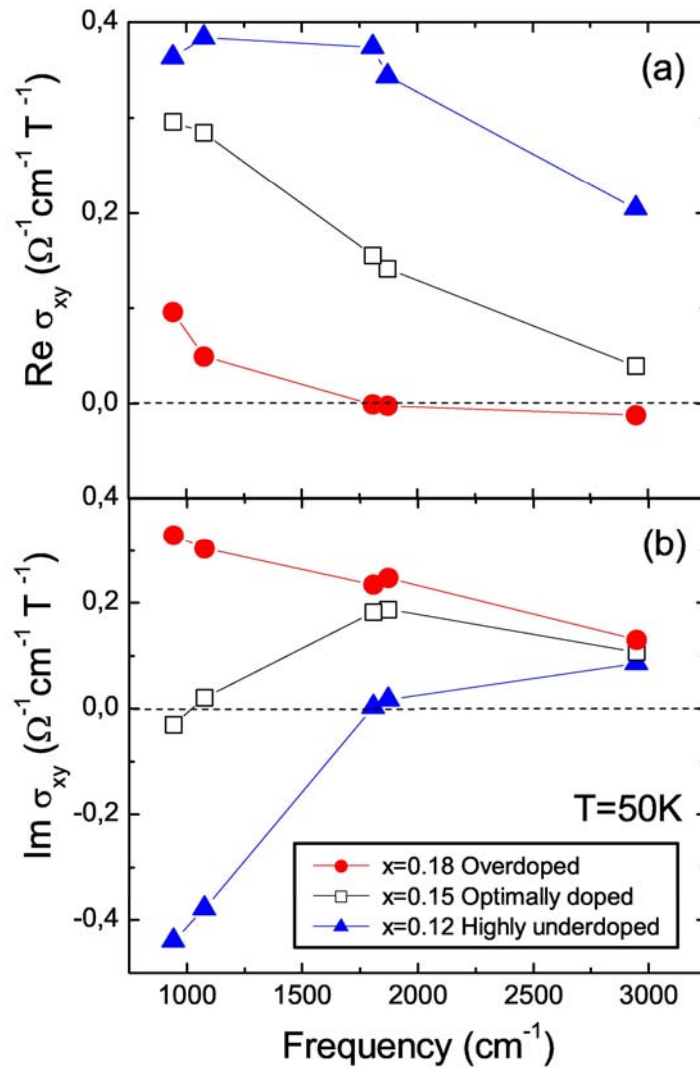
γ

Zimmers
IR data

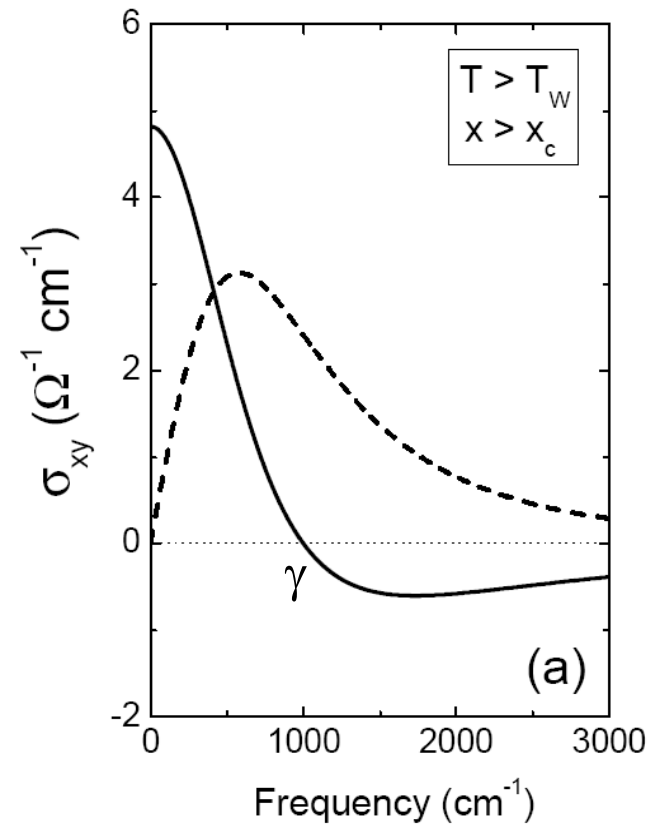


$1 + \lambda$

e-doped: IR Hall in $\text{Pr}_{1.82}\text{Ce}_{0.18}\text{CO}_4$



Drude model

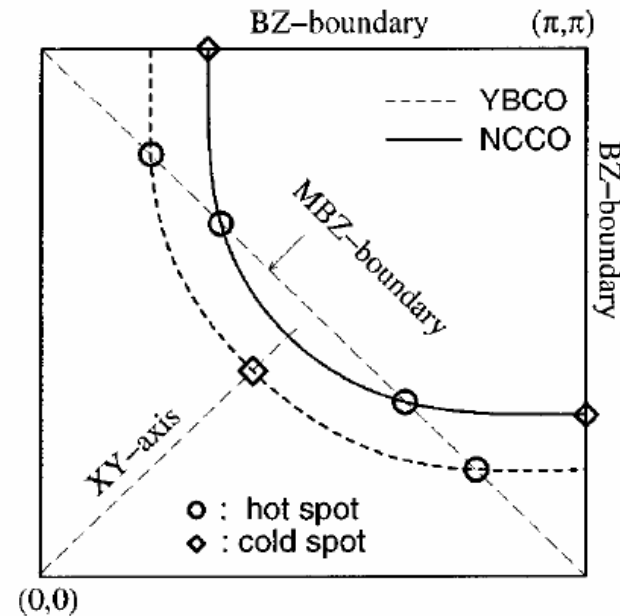


$$\sigma_{xy} = \frac{(ne^2 / m^*)(eB / m^* c)}{(\gamma - i\omega)^2}$$

Hall and longitudinal scattering

$\omega = 1000 \text{ cm}^{-1}$, $T = 100 \text{ K}$

	Bi2212 Opt.	PCCO X=0.18
γ_{xx}^* , cm^{-1}	750	750
γ_{xy}^* , cm^{-1}	360	1300



$$\gamma^* = \frac{\gamma}{1 + \lambda} = \frac{\gamma}{m^* / m_b}$$

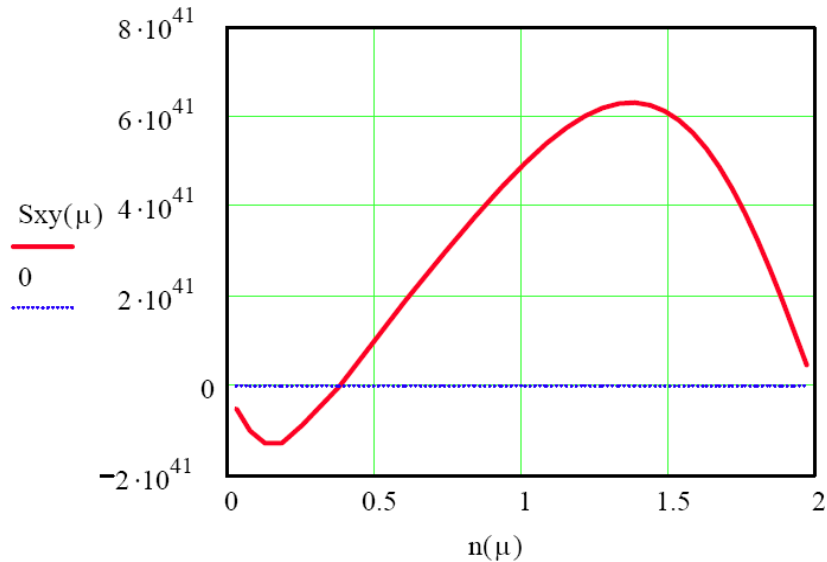
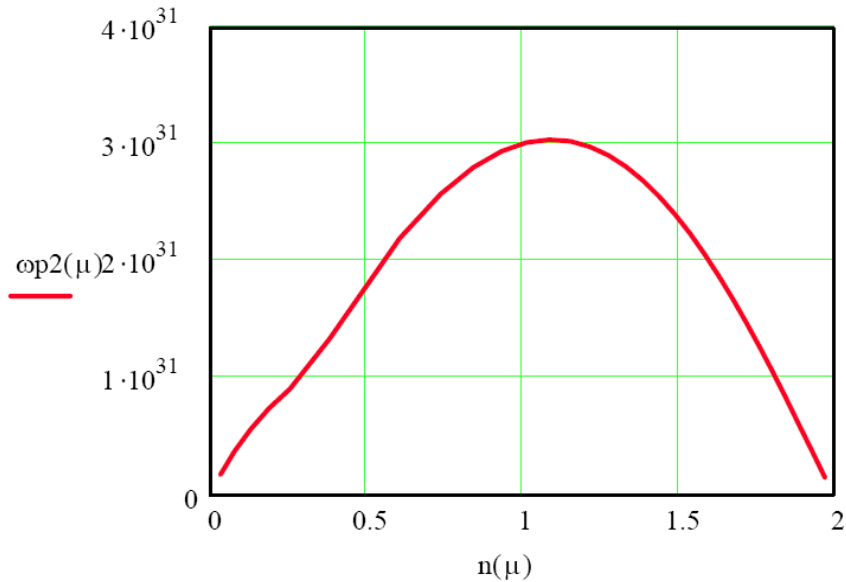
Schmadel et al., PRB (2007)

Zimmers et al., cond-mat/0510085

$$\sigma_{xy} = \frac{e^2}{hd} \frac{eB}{\hbar c} \oint \frac{dk \cdot v_k \times dv_k / dk}{(\gamma - i\omega)^2}$$

Conductivity Sum Rules: Cuprates

Band values



$$E(k_x, k_y) = -2t_1(\cos(k_x) + \cos(k_y)) + 4t_2 \cos(k_x) \cos(k_y) - 2t_3(\cos(2k_x) + \cos(2k_y))$$

$$t_1 = 0.38 \text{ eV}, t_2 = 0.32t_1, t_3 = 0.5t_2$$

$$S_{xx} \equiv \int_0^{\omega_c} \frac{2}{\pi} \text{Re}(\sigma_{xx}) d\omega \equiv \frac{e^2}{2c_0} \sum_k \text{Tr}(m_k^{-1}) n_k$$

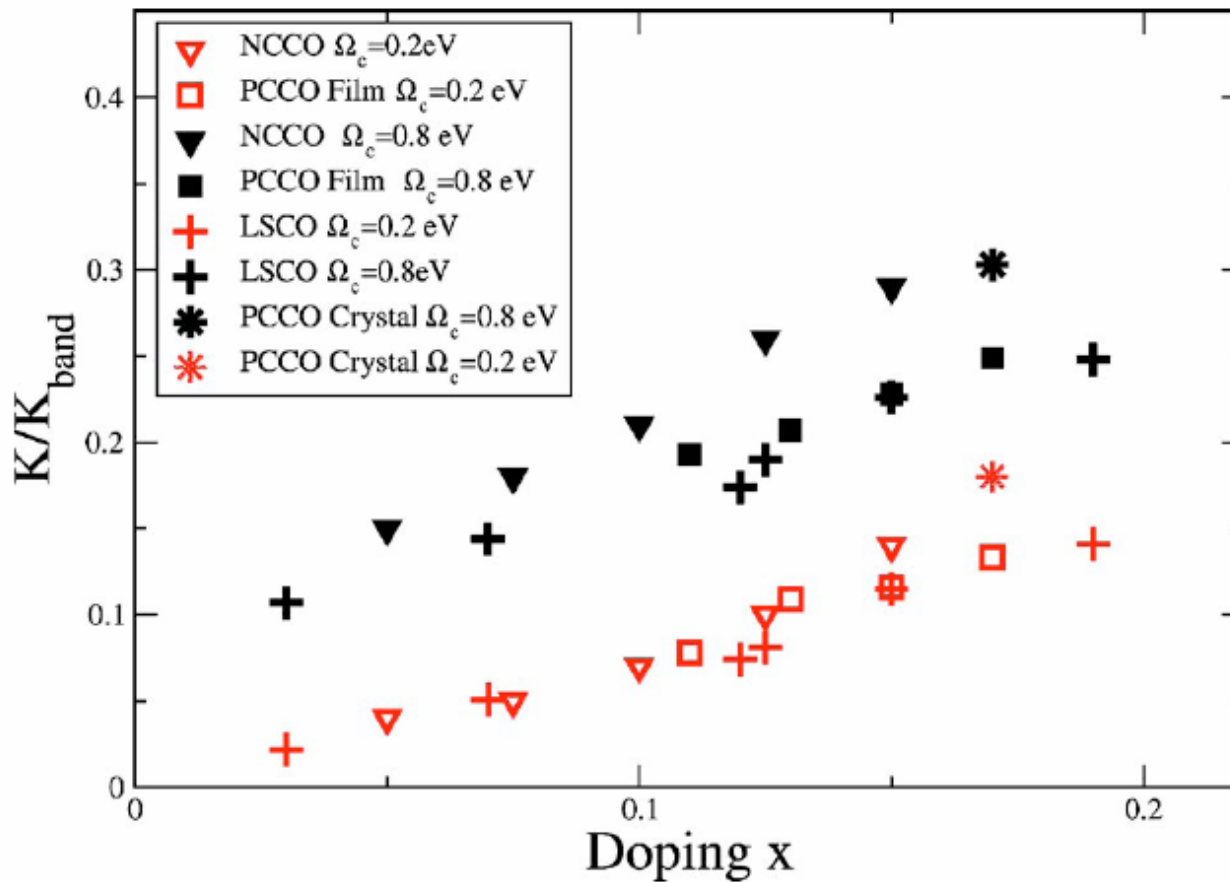
$$S_{xy} \equiv \int_0^{\omega_c} \frac{2}{\pi} \text{Im}(\sigma_{xy}) \omega d\omega \equiv \frac{e^3 B}{2c_0} \sum_k \det(m_k^{-1}) n_k$$

$$S_\theta \equiv \int_0^{\omega_c} \frac{2}{\pi} \text{Re}(\theta_H) d\omega = \frac{S_{xy}}{S_{xx}} \equiv \omega_H$$

σ_{xx} partial sum rule

LSCO, NCCO, PCCO

$$S_{xx} \equiv \int_0^{\Omega_c} \frac{2}{\pi} \text{Re}(\sigma) d\omega$$



Conductivity Sum Rules

$$S_{xx} \equiv \int_0^{\Omega_c} \frac{2}{\pi} \text{Re}(\sigma) d\omega$$

$$S_{xx}^{band} = \frac{e^2}{2} \sum_k \text{Tr}(m_k^{-1}) n_k \cong \frac{ne^2}{m^*}$$

$$S_{xy} \equiv \int_0^{\Omega_c} \frac{2}{\pi} \text{Im}(\sigma_H) \omega d\omega$$

$$S_{xy}^{band} = e^3 B \sum_k \det(m_k^{-1}) n_k \cong \frac{ne^2}{m^*} \frac{eB}{m^* c}$$

Bi2212 opt. doped

$$\frac{S_{xx}^{exp}}{S_{xx}^{band}} \approx 0.33$$

$$\frac{S_{xy}^{exp}}{S_{xy}^{band}} \approx 0.09$$

Schmadel et al.,
PRB (2007)

PCCO x=0.18

$$\frac{S_{xx}^{exp}}{S_{xx}^{band}} \approx 0.25$$

$$\frac{S_{xy}^{exp}}{S_{xy}^{band}} \approx 0.09$$

Zimmers et al.,
cond-mat/0510085

Conclusions on Optimally doped Cuprates

1. $\gamma_{xy} \neq \gamma_{xx}$ vertex corrections

2. $\gamma_{xy}^{PCCO} > \gamma_{xx} > \gamma_{xy}^{Bi2212}$ Hot spots

3. $\frac{S_{xy}}{S_{xy}^{band}} < \frac{S_{xx}}{S_{xx}^{band}} < 1$ Mott physics:
e* or m*?

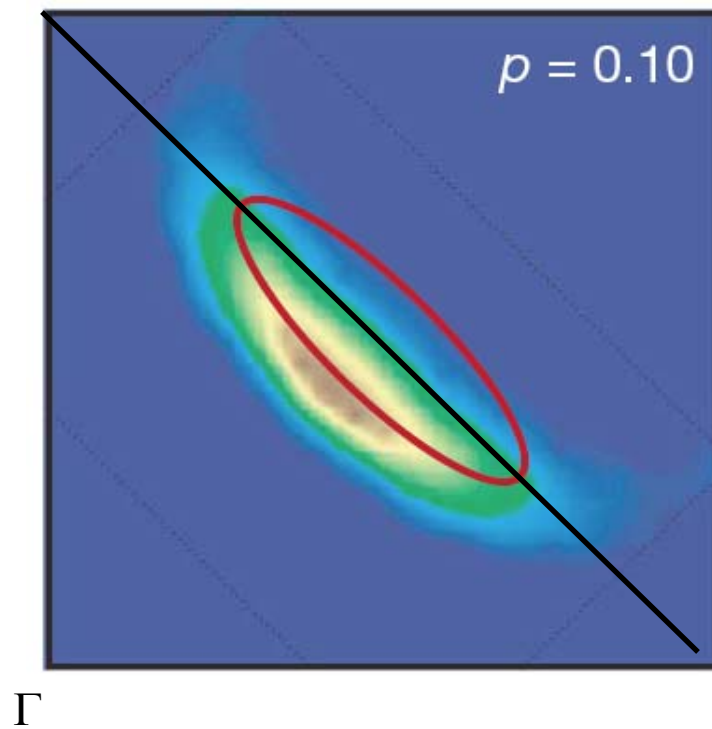
$$S_{xx} \propto \frac{e^2}{m}, \quad S_{xy} \propto \frac{e^3}{m^2}$$

IR Hall Effect
in
Underdoped Cuprates

ARPES Fermi Surface

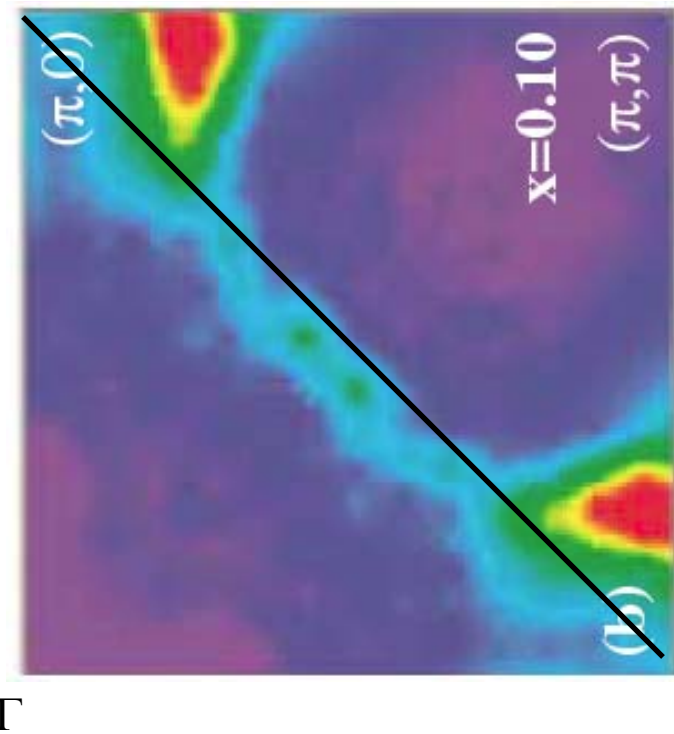
Hole doped

(π, π)



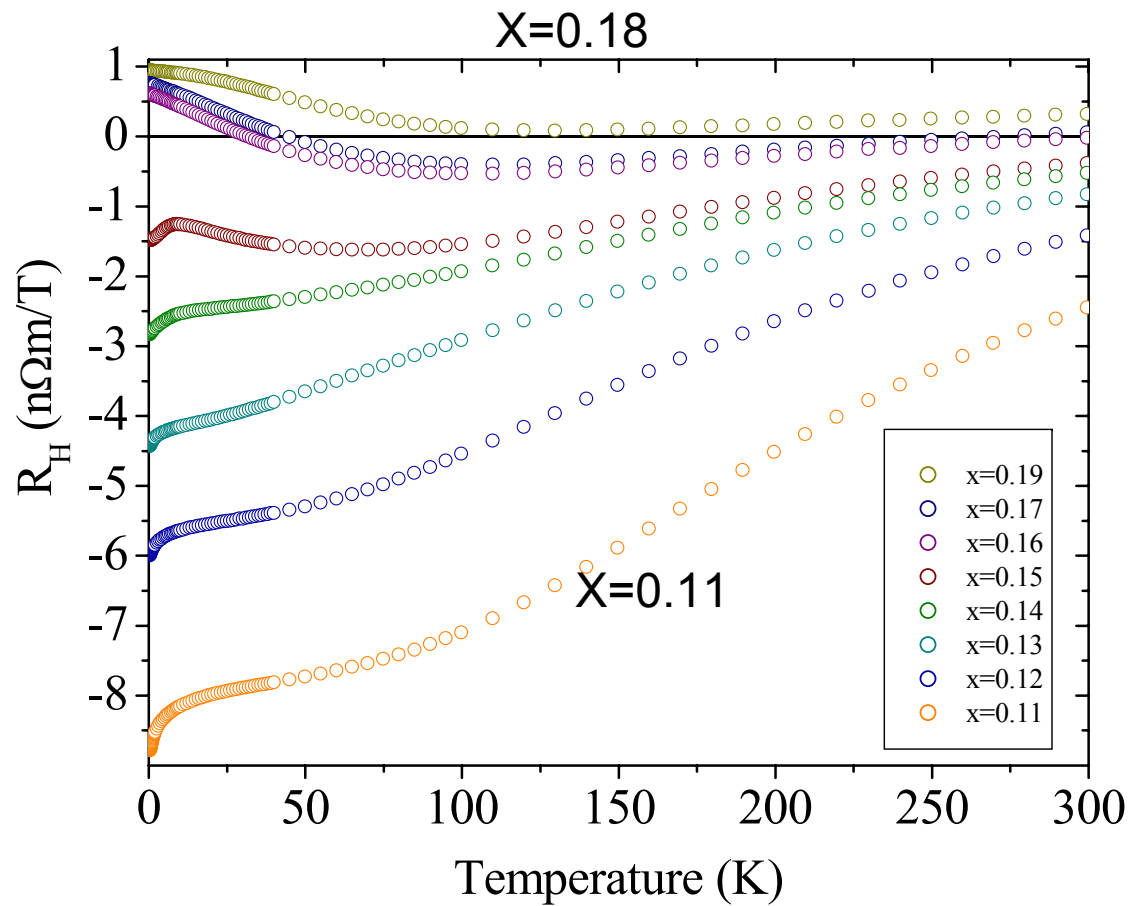
K. Shen, Science (2005)

Electron doped



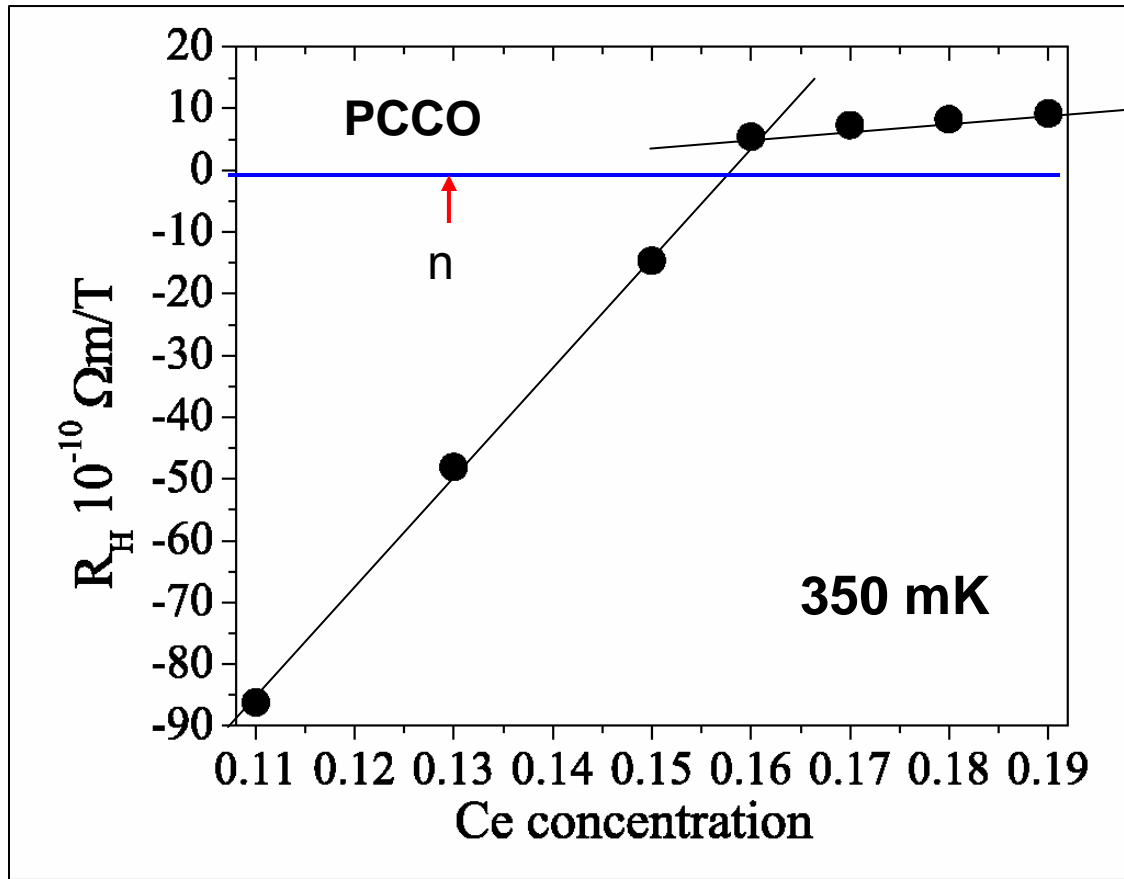
Armitage *et al.*, PRL (2002)

$\text{Pr}_{2-x}\text{Ce}_x\text{CuO}_4$ DC R_H



Y. Dagan *et al.*, *Phy. Rev. Lett.* **92**, 167001 (2004).

Hall coefficient $\text{Pr}_{2-x}\text{Ce}_x\text{CuO}_4$

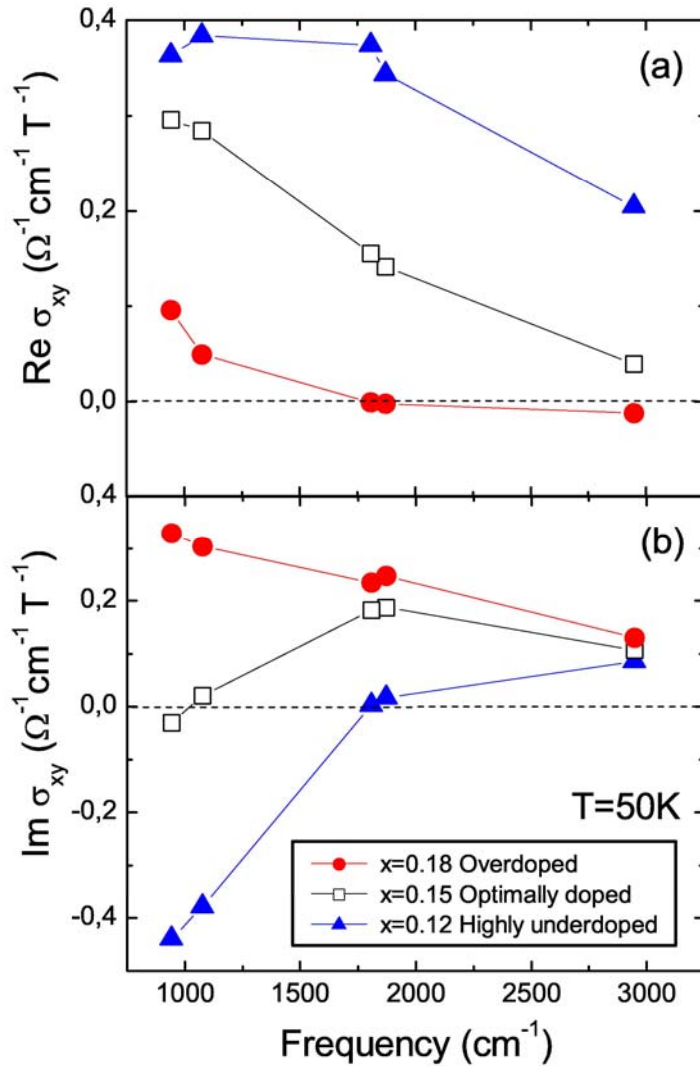


Quantum
critical
point?

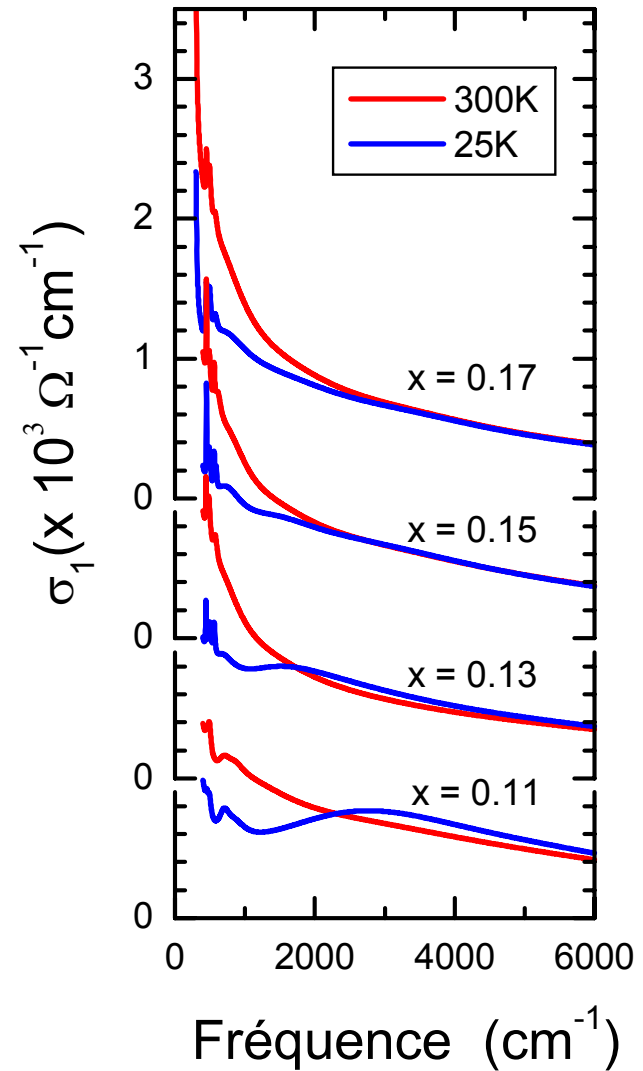
Hall - Dagan et al., PRL 92 (2004)

Neutrons - Motoyama, et al., Nature (2007)

e-doped: IR Hall in $\text{Pr}_{2-x}\text{Ce}_x\text{CO}_4$

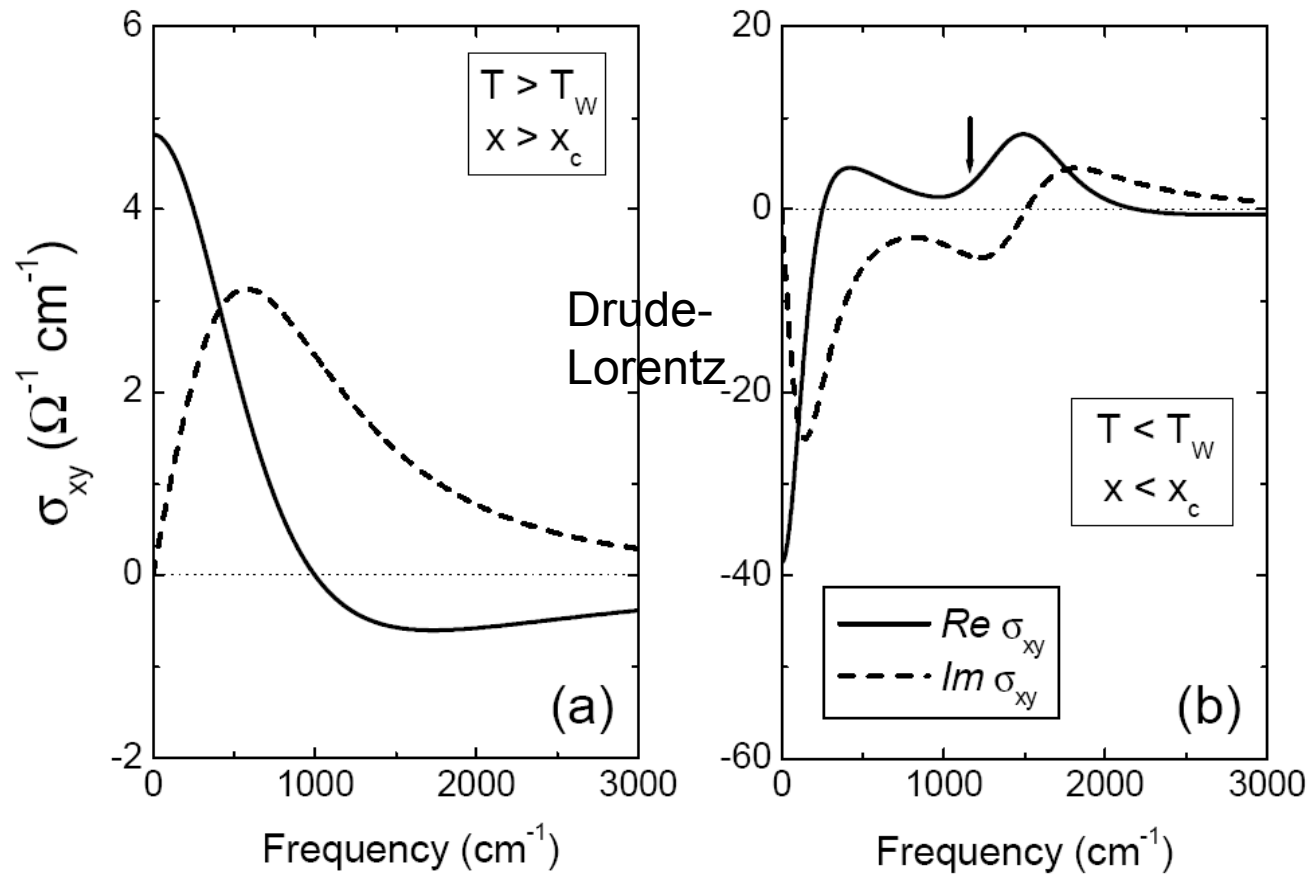


A. Zimmers et al., cond-mat/0510085



A. Zimmers et al., Europhysics Lett. (2005)

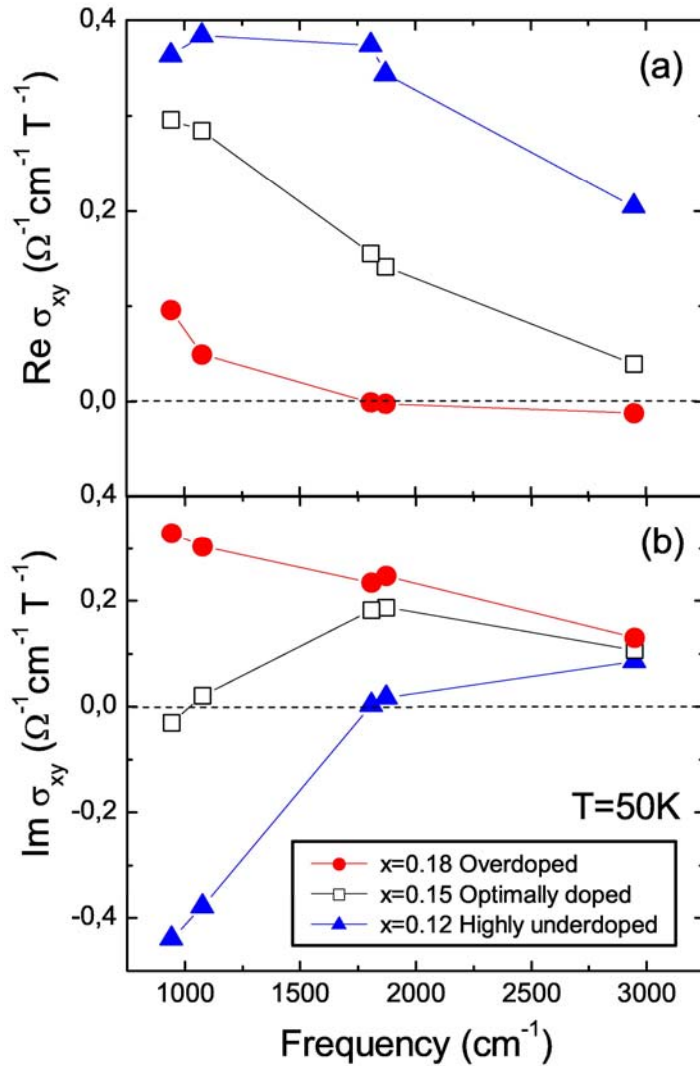
σ_{xy} in SDW: Drude-Lorentz model



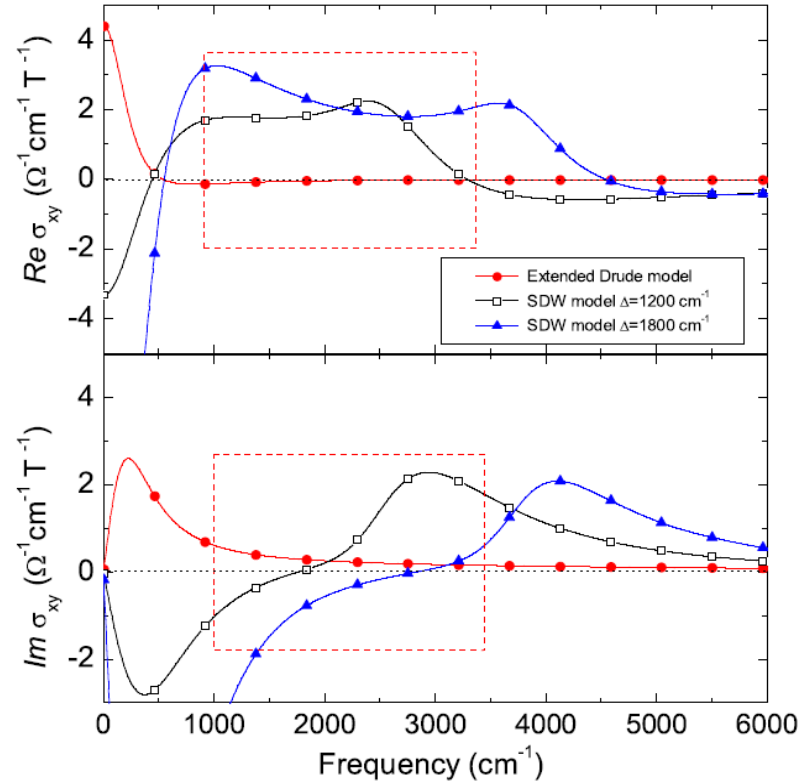
Drude

$$\sigma_{xy} = \frac{(ne^2 / m^*)(eB / m^* c)}{(\gamma - i\omega)^2}$$

Mid IR σ_H for PCCO vs. x



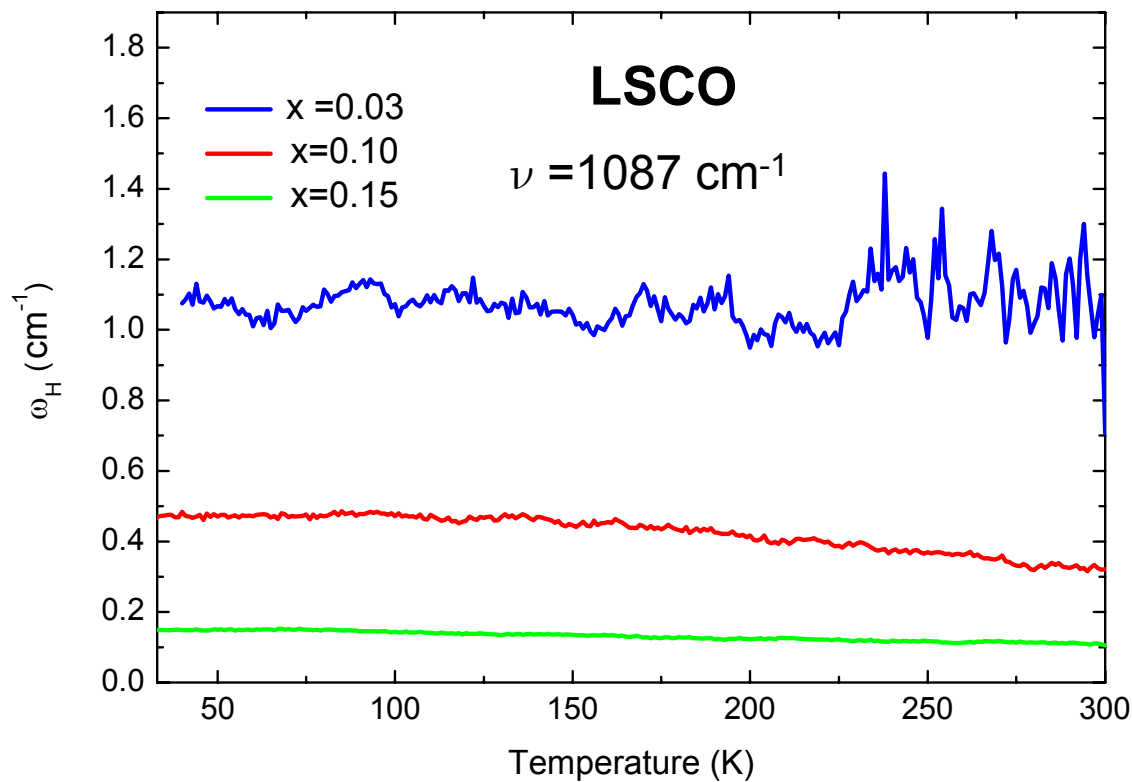
SDW model: J. Lin and A. Millis



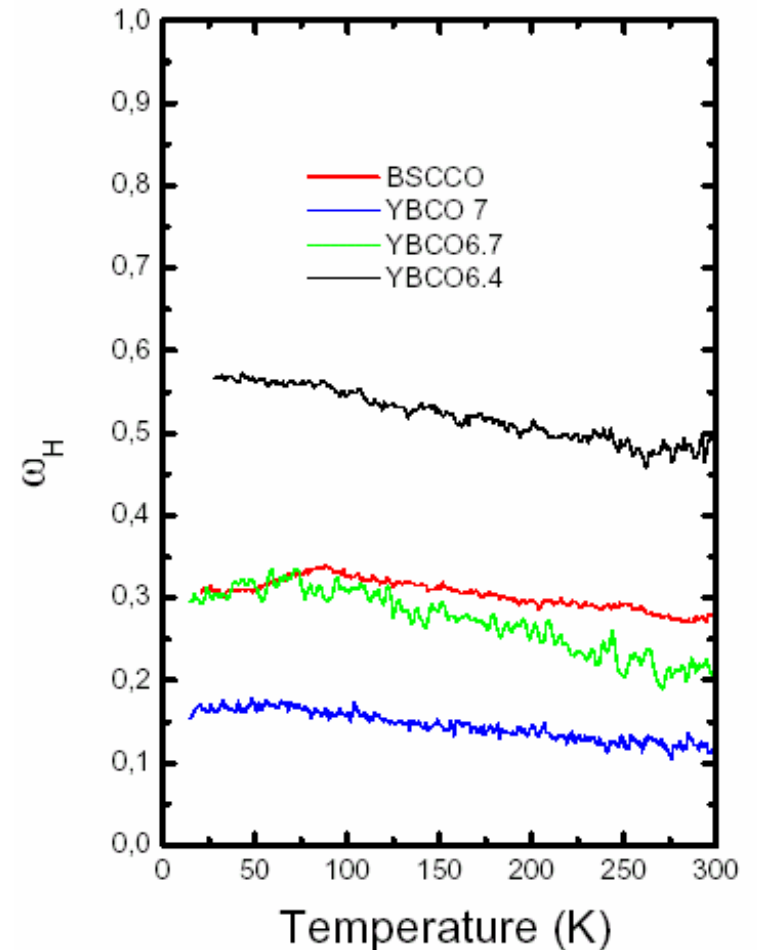
x	$2\Delta(\text{meV})$
0.11	450
0.15	300
0.18	0

Underdoped h-cuprates: Hall Frequency: ω_H

$$\theta_H^{-1} = \frac{\gamma_H - i\omega}{\omega_H}, \quad \omega_H \equiv \frac{eB}{m_H c}$$



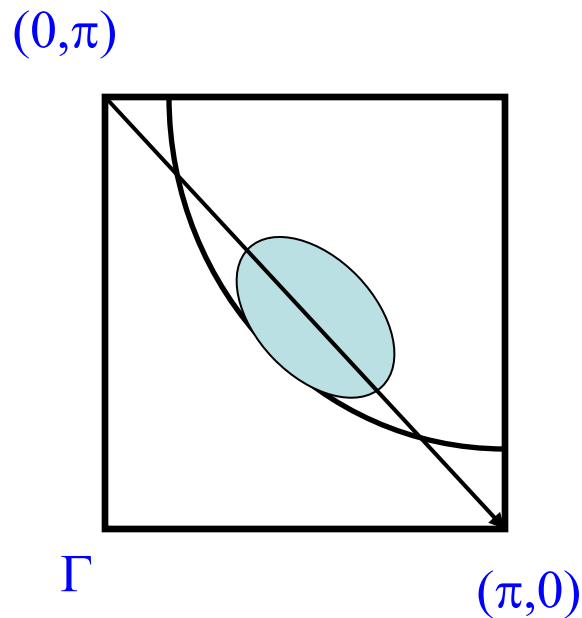
L. Shi, et al., cond-mat/0510794



L. Rigal et al., PRL (04)

Fermi Pocket Hall Frequency

Partially gapped Fermi surface



$$\theta_H = \frac{\sigma_{xy}}{\sigma_{xx}} \simeq \frac{\omega_H}{\gamma - i\omega} \quad \omega_H \equiv \frac{eB}{m_H c}$$

$$\omega_H = \frac{eB}{\hbar c^2} \frac{\oint dk \cdot v \times dv/dk}{\oint dk |v|} = \frac{eB}{\hbar c^2} \frac{\oint dv \times v}{\oint dk |v|}$$

Chakravarty et al., Phys. Rev. (2004).

Konatani theory also predicts increase in ω_H

Taillefer et al., Nature (2007)

Kanigel et al., Nature Physics, (2006)

Conclusions

Underdoped cuprates

1. e-doped: SDW gap evidence in σ_{xx} and σ_{xy}
 $2\Delta \sim 300 \text{ meV} - 450 \text{ meV}$
2. h-doped: evidence for Fermi surface reconstruction:
if gap $\rightarrow 2\Delta > 500 \text{ meV}$
3. Is the reconstructed FS real, fluctuating or only effective?
4. Hall conductivity spectroscopy gives important information on strongly interacting electron materials