

# Carbon Nanotubes as Nanoelectromechanical Systems

- Introduction: NEMS
- Model for suspended carbon nanotubes
- Displacement and eigenmodes
- Stability diagram
- Transport (preliminary)
- Outlook and conclusions

## Quantum Transport

- Pablo Jarillo-Herrero
- Sami Sapmaz
- Chris Lodewijk
- Leonid Gurevich
- Herre van der Zant
- Leo Kouwenhoven

## Theoretical Physics

- Omar Usmani
- Wouter Wetzels
- Wataru Izumida
- Milena Grifoni
- Y.M.B.

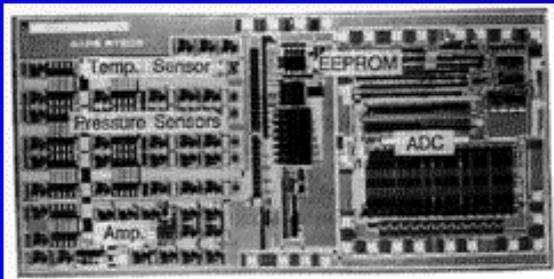
## Molecular biophysics

- Jing Kong
- Abdou Hassanien
- Cees Dekker

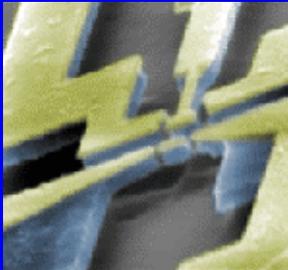
# Nanoelectromechanical systems



Electromechanical systems



Microelectromechanical systems



Nanoelectromechanical systems

# Nanoelectromechanical systems

**NEMS** – nanoscale devices which convert electrical current into mechanical energy or vice versa.

## Experiments: precise measurements

attoNewtons of force (*Stowe et al '97*)

electrometry (*Cleland and Roukes '98*)

quantum of thermal conductance (*Schwab et al '00*)

Casimir force (*Chan et al '01*)

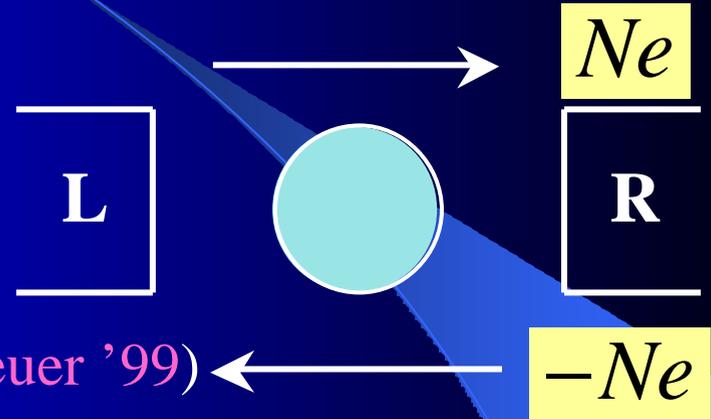
**Possible applications:** nanoscale sensors and actuators

# NEMS research directions

➤ **Shuttling:** First theoretical proposal by Gorelik *et al* '98

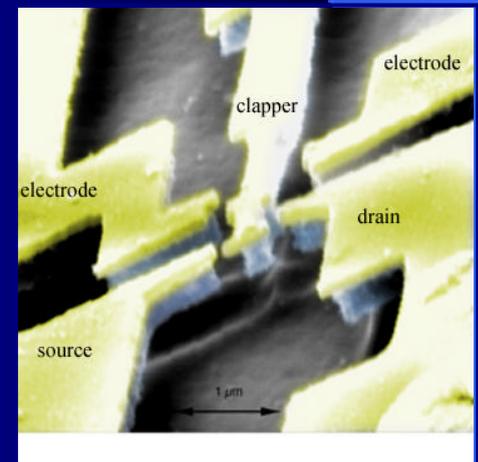
## Experiments:

- Classical shuttle (Erbe *et al* '98)
- Silver grain (Tuominen, Krotkov, Breuer '99)
- Fullerene molecule (Park *et al* '00)
- ac driven cantilever (Erbe *et al* '01)



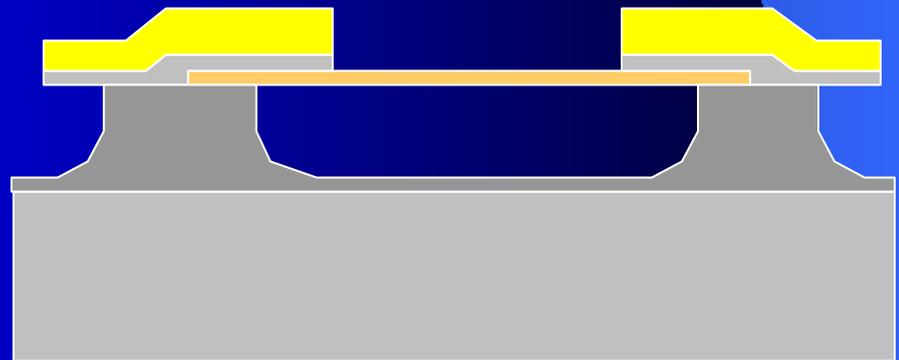
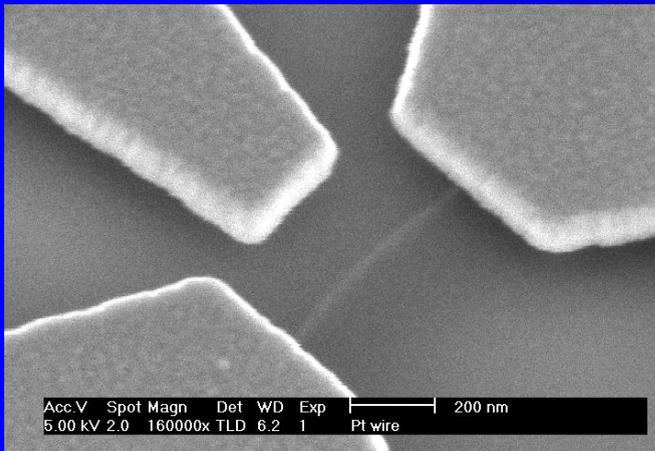
➤ **Suspended beams**

- Electromechanical noise
- Bistability and quantum effects
- Position and eigenmodes



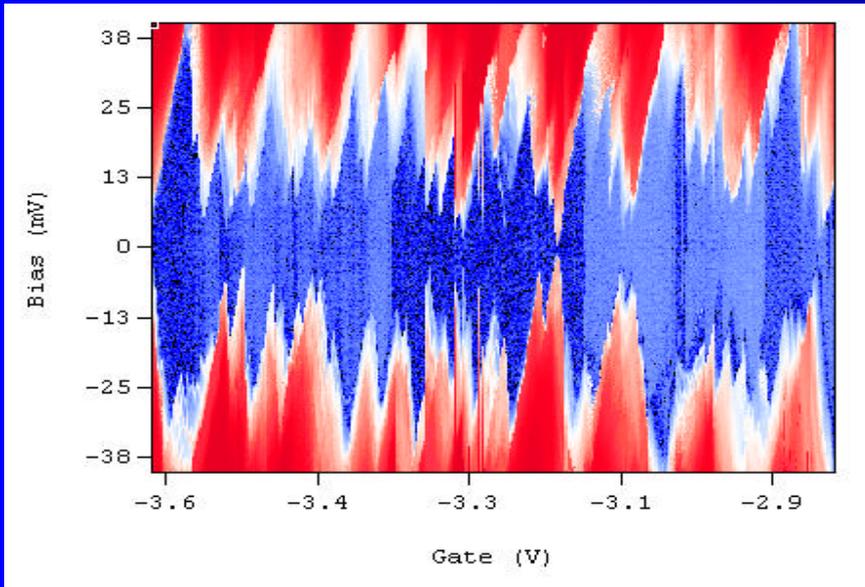
# Suspended beams - experiments

- ❖ Silicon quantum dot embedded into a suspended beam  
(Höhberger *et al* '02)
- ❖ Suspended doubly-clamped carbon nanotubes  
(P. Jarillo-Herrero *et al*, in progress)

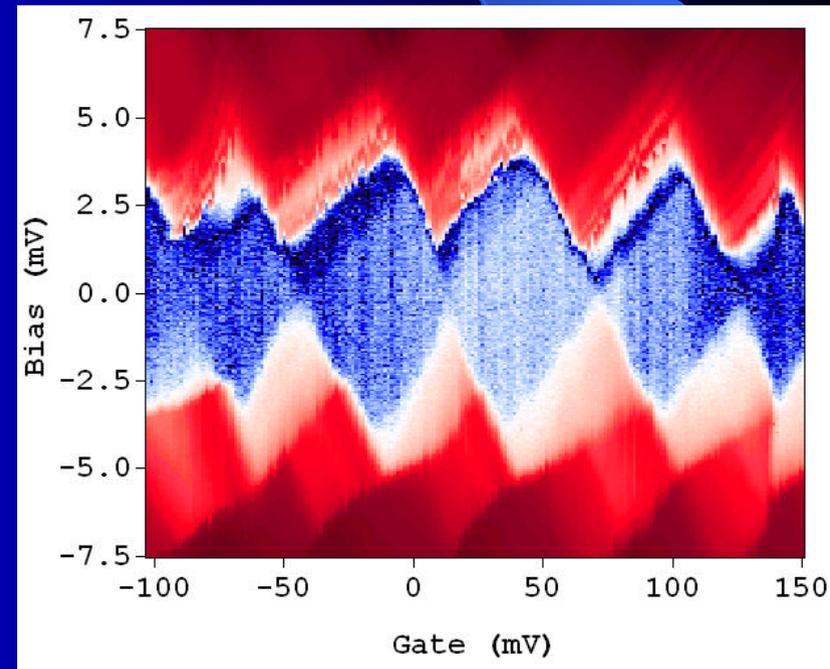


# Stability diagram – Delft experiments

$L=1200\text{nm}$ ,  $T=300\text{mK}$



Before under-etching

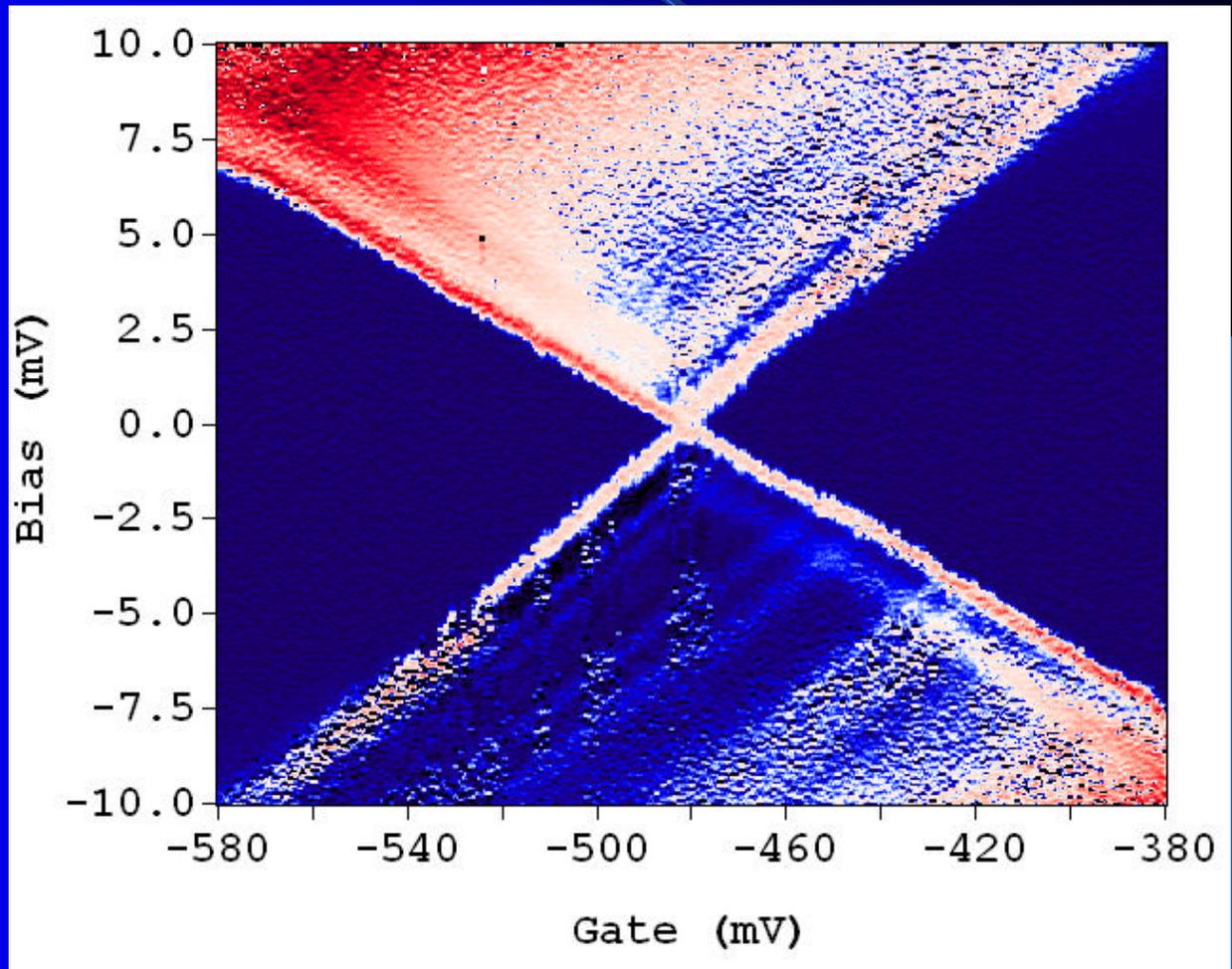


After under-etching



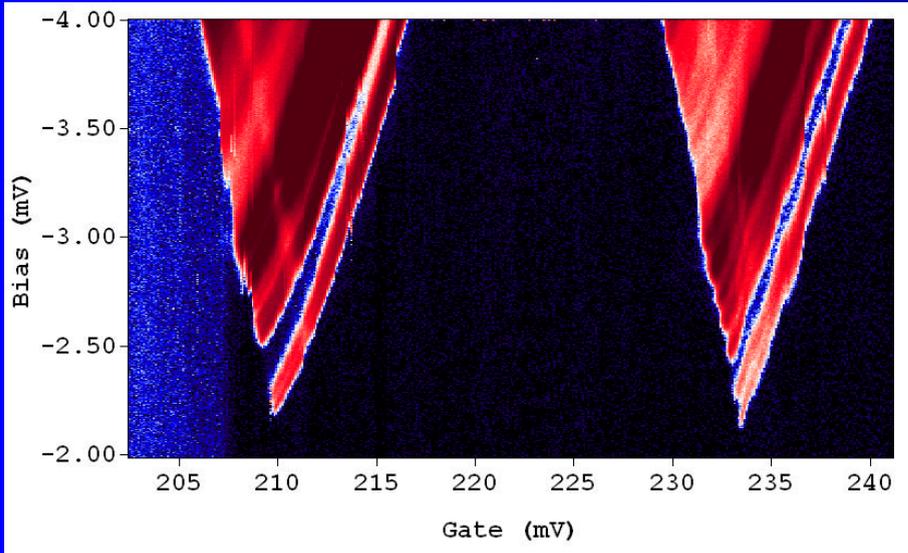
# Stability diagram – Delft experiments

P. Jarillo-Herrero *et al*, in progress  $L=140\text{nm}$ ,  $T=300\text{mK}$



# Stability diagram – Delft experiments

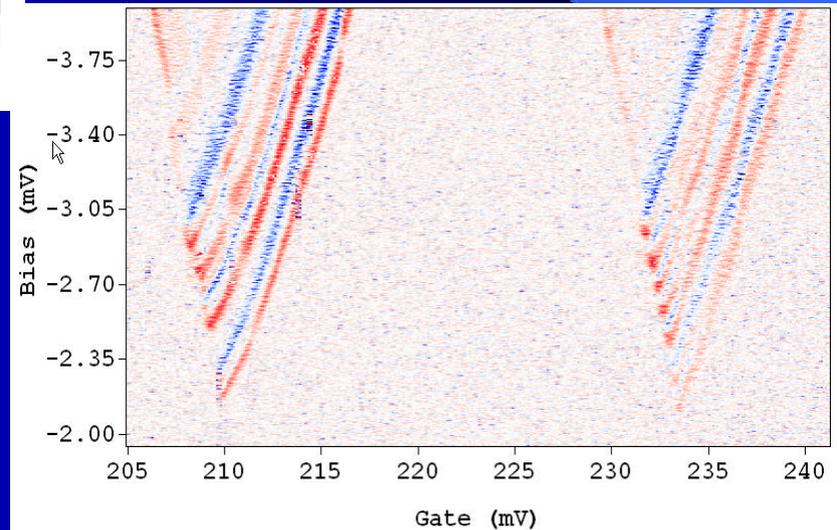
P. Jarillo-Herrero *et al*, in progress



$\ln I$

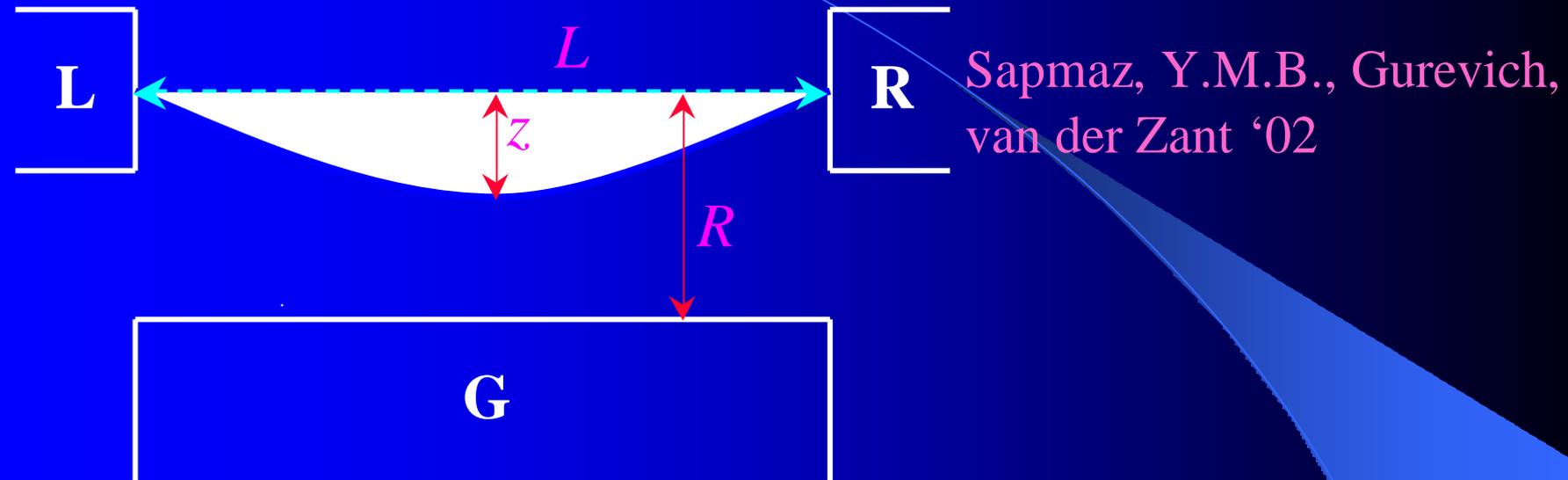


$dI / dV$



Inelastic tunneling?

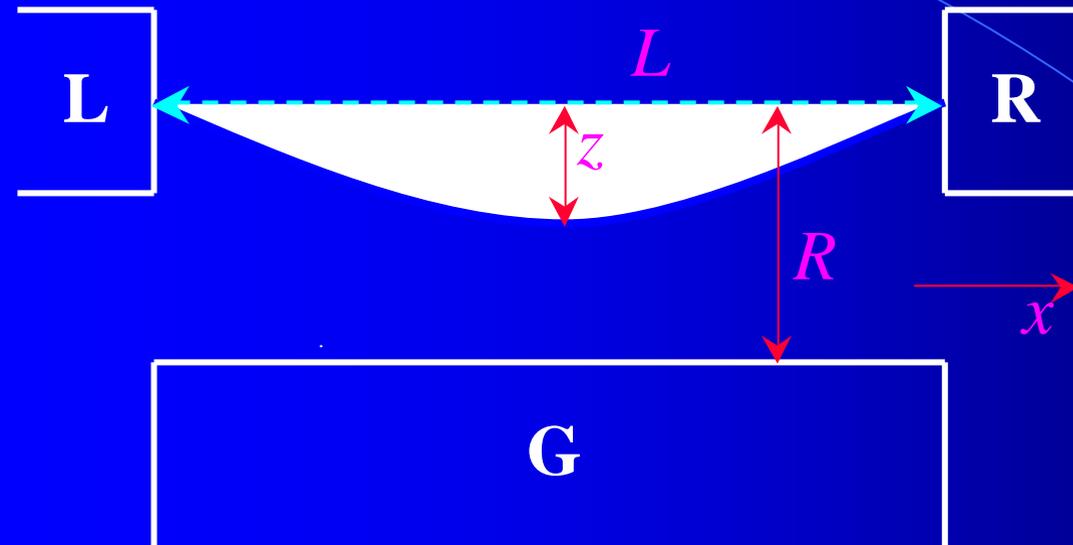
# Modeling suspended nanotubes



## Features of the model:

- Interaction effects taken into account via charging energy;
- Mechanical degrees of freedom via classical theory of elasticity; nanotube modeled as an elastic rod

# Elastic energy



Elastic modulus

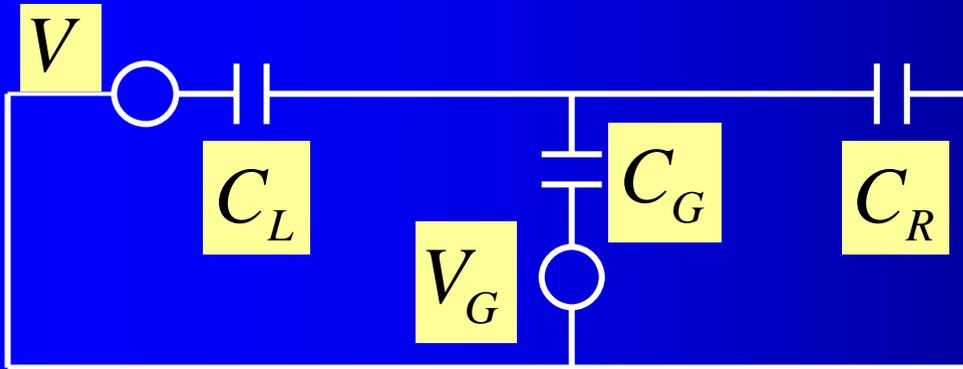
$$W_{el}\{z(x)\} = \int_0^L dx \left\{ \frac{EI}{2} z''^2 + \frac{T}{2} z'^2 \right\}$$

Residual stress

$$T = T_0 + \frac{ES}{2L} \int z'^2 dx$$

Stress induced by bending

# Electrostatic energy



$$C_L, C_R \square C_G$$

$$C_G = \int \frac{dx}{2 \ln \frac{2(R - z[x])}{r}}$$

$$W_{e-st} \{z[x]\} = \frac{(ne)^2}{2C_G \{z\}} - neV_G \approx \frac{(ne)^2}{2C_0} - neV_G - \frac{(ne)^2}{L^2 R} \int z[x] dx$$

Electrostatic force

# Displacement

Equations of motion:

$$IEz'''' - Tz'' = K_0 \equiv \frac{(ne)^2}{L^2 R}$$

1. To be solved for the displacement;
2. Stress to be found self-consistently

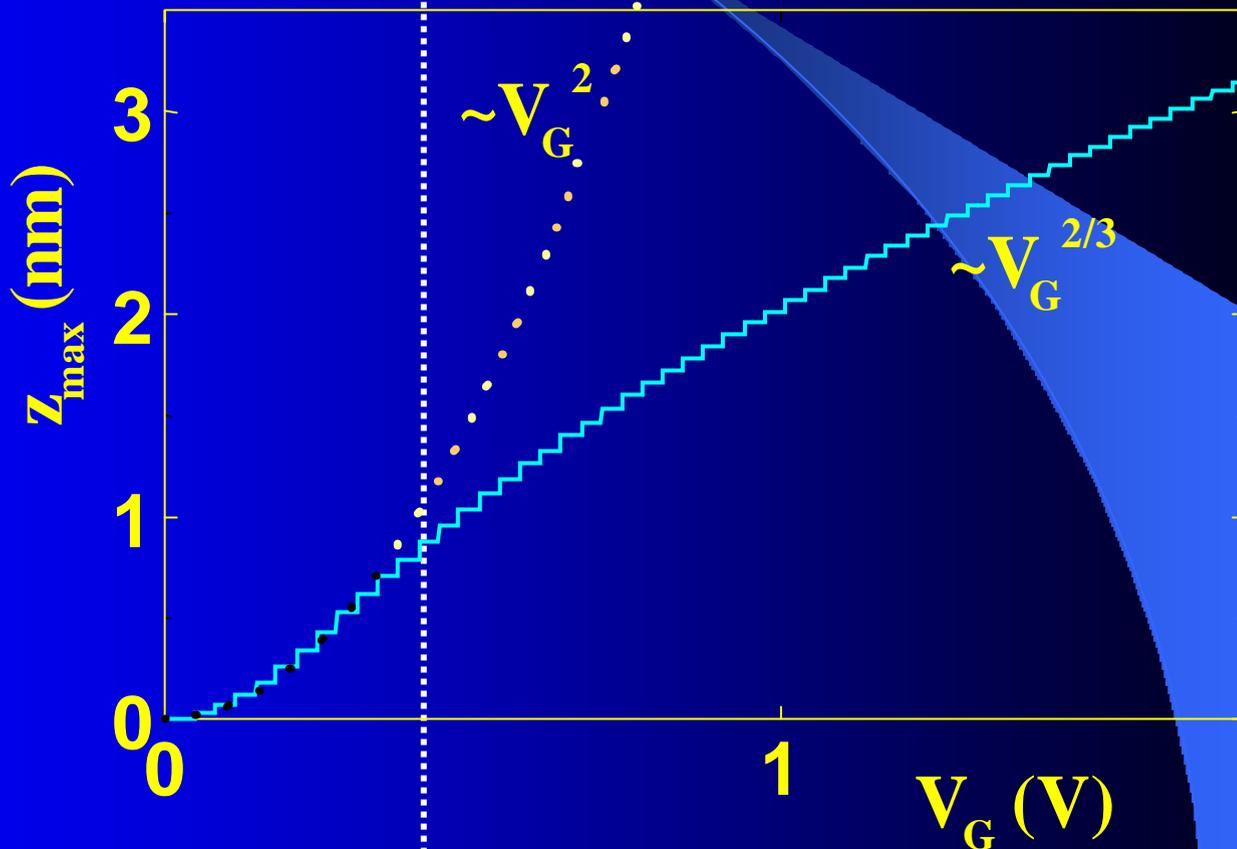
Results: Two regimes

•Weak bending:  $z_{\max} < r, z_{\max} \propto V_G^2$

•Strong bending:  $z_{\max} > r, z_{\max} \propto V_G^{2/3}$

# Displacement

$L = 500\text{nm}$   
 $R = 100\text{nm}$   
 $C_L = C_R = 0$



Weak bending



Strong bending

# Eigenmodes

Equation: 
$$\cosh y_1 \cos y_2 - \frac{1}{2} \frac{y_1^2 - y_2^2}{y_1 y_2} \sinh y_1 \sin y_2 = 1$$

$$y_{1,2} = \frac{L}{\sqrt{2}} \sqrt{\sqrt{\mathbf{x}^2 + \mathbf{l}^2} \pm \mathbf{x}^2}; \mathbf{x} = \sqrt{\frac{T}{EI}}; \mathbf{l} = \sqrt{\frac{4rS}{EI}} \mathbf{w}$$

Results for the fundamental mode

Weak bending: 
$$\mathbf{w}_0 = \sqrt{\frac{EI}{rS}} \frac{22.38}{L^2} + (\dots) V_G^4$$

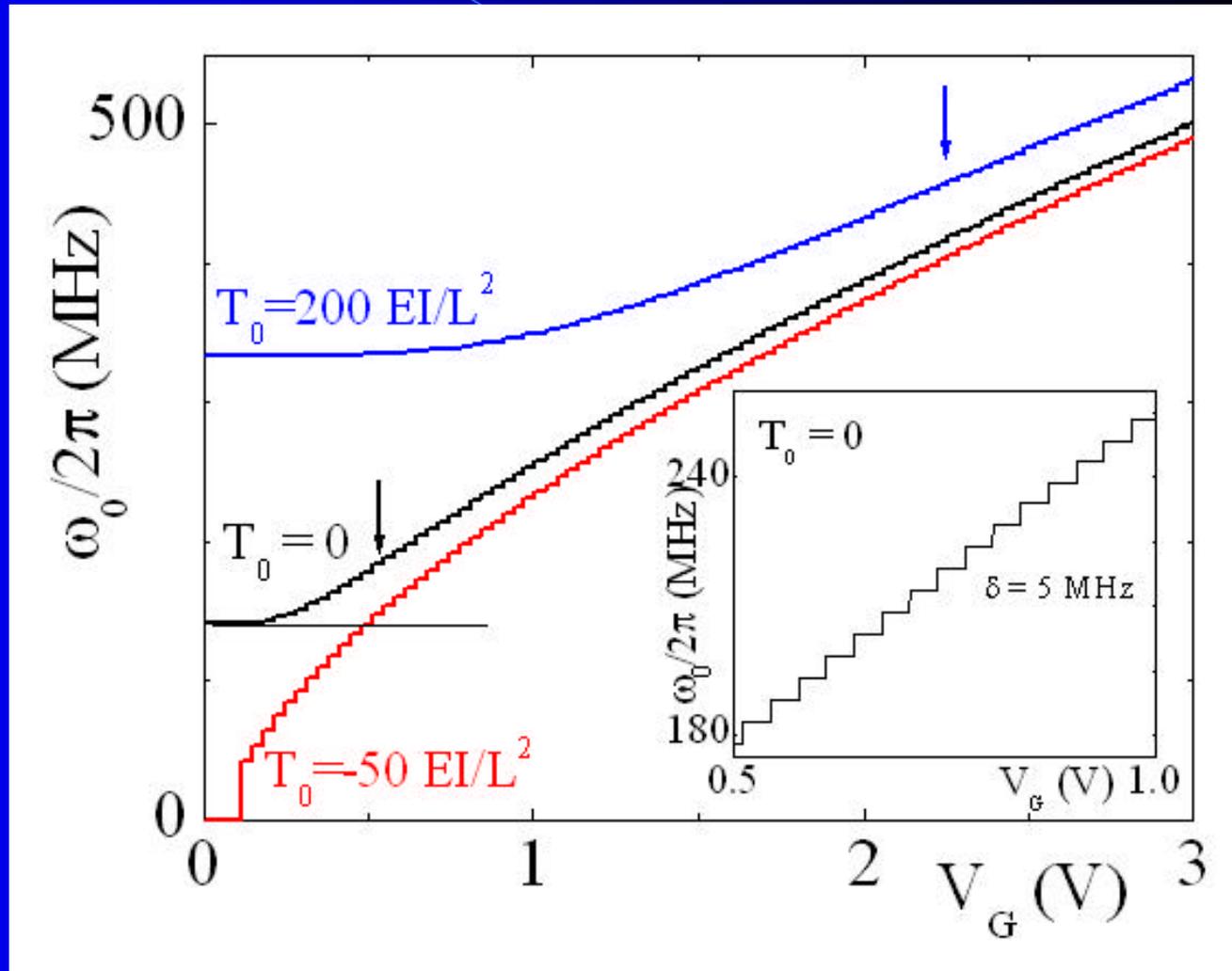
Strong bending: 
$$\mathbf{w}_0 \propto V_G^{2/3}$$

# Eigenmodes

$$L = 500\text{nm}$$

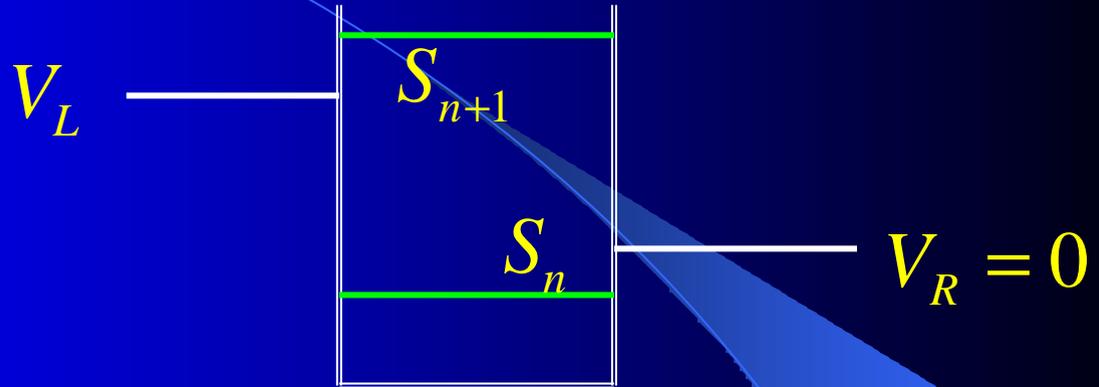
$$R = 100\text{nm}$$

$$C_L = C_R = 0$$



# Stability diagram for quantum dots

$$S_n = W_{n+1} - W_n$$
$$\approx \left(n + \frac{1}{2}\right) \frac{e^2}{C_G} - eV_G$$



Conditions that current is **not** flowing:

(a)  $S_{n+1} > eV_L$

(b)  $S_n < eV_L$

(c)  $S_{n+1} > 0$

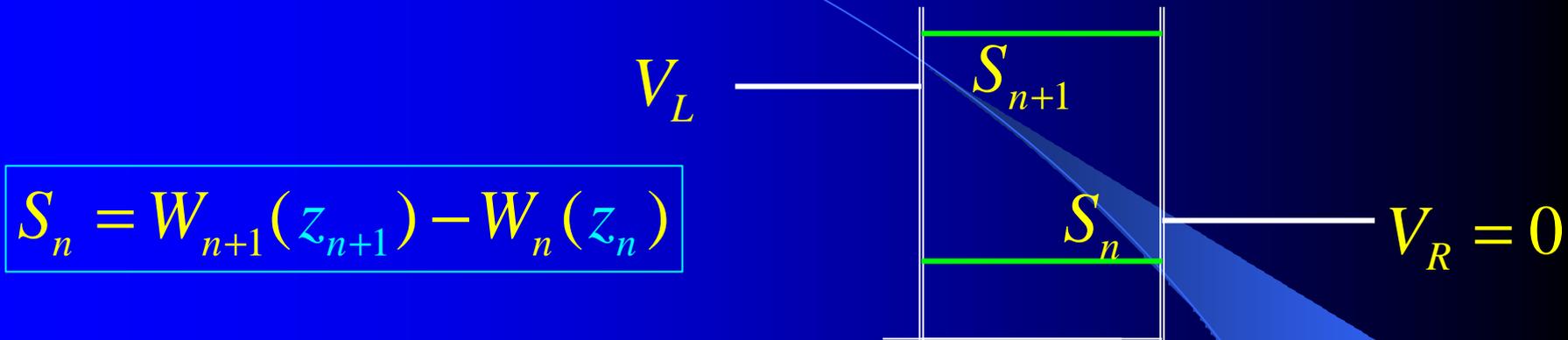
(d)  $S_n < 0$

Linear dependence



Coulomb diamonds

# Stability diagram for suspended nanotubes



Conditions that current is not flowing: Non-linear

Consequences:

- “Curvilinear diamonds”
- Diamonds become smaller with increasing the charge

Magnitude of the effect: small

Probably not observable experimentally

# Transport

Suspended nanotubes: damped oscillators

Quality factor:  $Q = \gamma / \omega_0 \approx 10^2 \div 10^3$  (Reulet et al '00)

Weakly damped oscillators!

“Stationary regime”: an external force induces oscillations with the frequency of the fundamental mode.

$$q = W \sin \omega_0 t$$

# Transport

Usmani & Y.M.B., in preparation

Two charge states:  $n=0$  &  $n=1$

Time-dependent tunnel rates

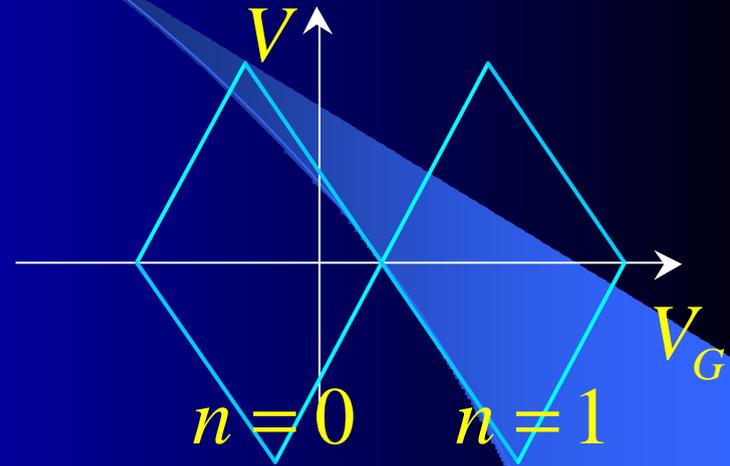
Time-dependent occupation probabilities

Stochastic force due to switching between the states

Energy dissipation over the period

Amplitude found self-cons.

$$C_L = C_R = 0$$



Current

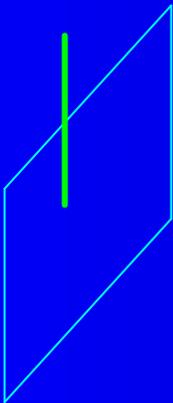
# Tunnel rates

$$C_L = C_R = 0$$

$$\Gamma_L^{0 \rightarrow 1}(t) = \frac{1}{e^2 R} \left[ \frac{e^2}{2C_G} + eV_G - \mathbf{a}W \sin \mathbf{w}_0 t \right] \mathbf{q} \left[ \frac{e^2}{2C_G} + eV_G - \mathbf{a}W \sin \mathbf{w}_0 t \right]$$

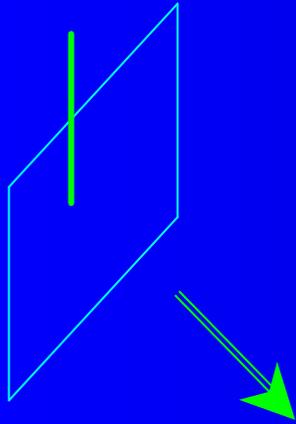
$$\Gamma_R^{1 \rightarrow 0}(t) = \frac{1}{e^2 R} \left[ -\frac{e^2}{2C_G} + e(V - V_G) + \mathbf{a}W \sin \mathbf{w}_0 t \right] \mathbf{q} \left[ -\frac{e^2}{2C_G} + e(V - V_G) + \mathbf{a}W \sin \mathbf{w}_0 t \right]$$

$$\Gamma_L^{1 \rightarrow 0}(t) = \Gamma_R^{0 \rightarrow 1}(t) = 0$$



$$\mathbf{a} = \frac{e^2}{LR} \quad \text{- force for } n=1$$

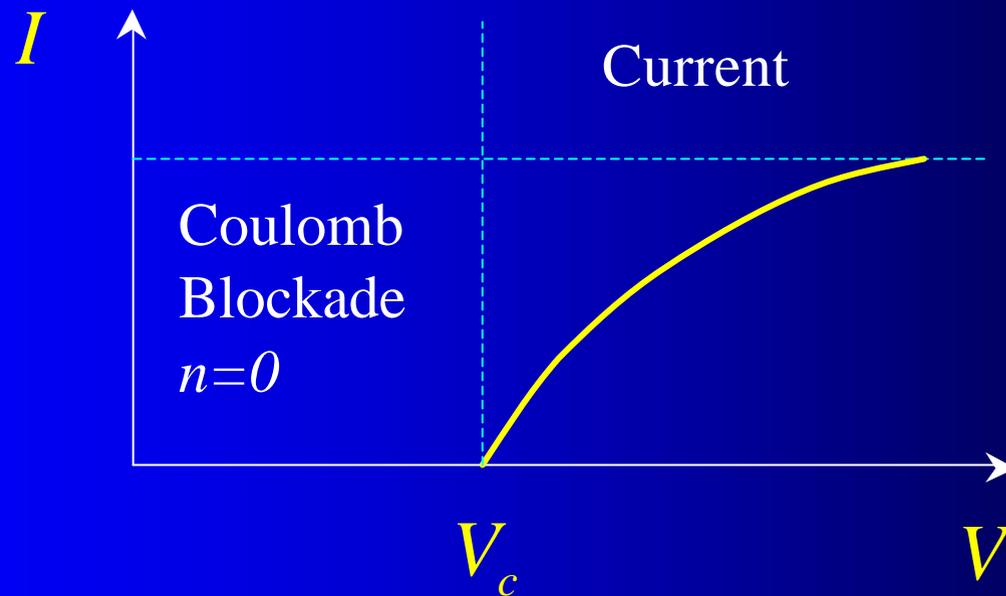
# Tunnel rates



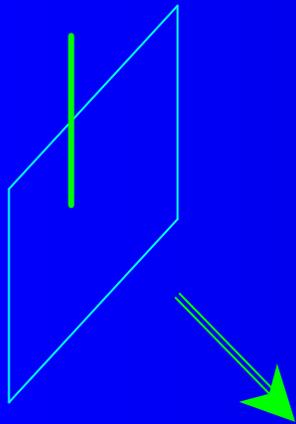
$$a = \frac{e^2}{LR}$$

- force for  $n=1$

Without mechanical oscillations:

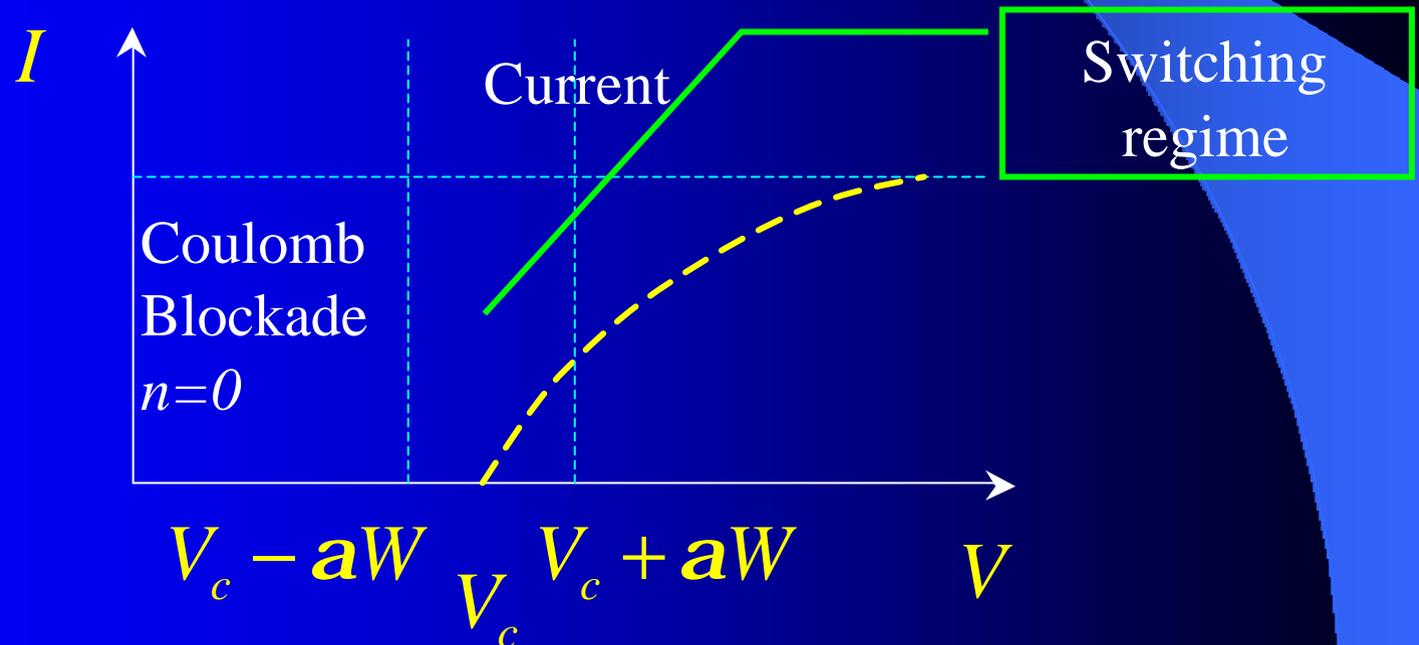


# Tunnel rates



$$a = \frac{e^2}{LR} \quad \text{- force for } n=1$$

With mechanical oscillations:



# Occupation probabilities

$P_1(t)$  - probability to have charge  $e$

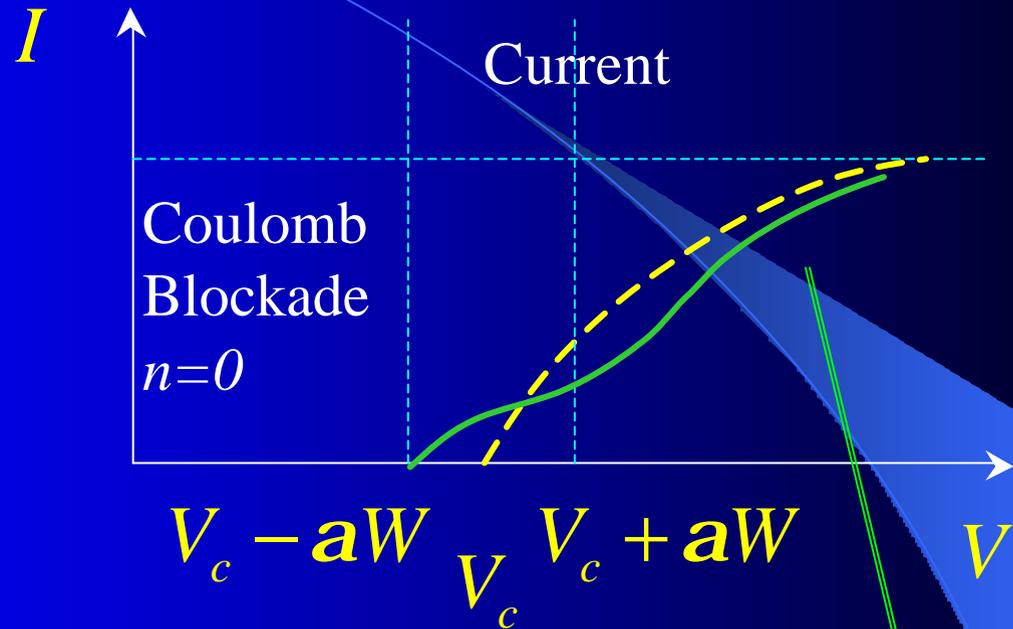
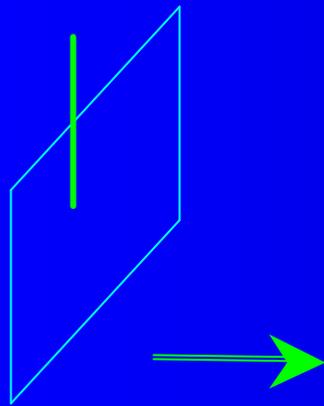
$P_0(t) = 1 - P_1(t)$  - probability to have charge  $0$

$$\frac{dP_1}{dt} = \Gamma_L^{0 \rightarrow 1}(t) P_0(t) - \Gamma_R^{1 \rightarrow 0}(t) P_1(t)$$

Current:

$$I = \frac{e}{2} \left[ \Gamma_R^{1 \rightarrow 0} P_1 + \Gamma_L^{0 \rightarrow 1} P_0 \right]$$

# Current



$$dI \propto -\frac{Q}{w_0^3} \left( \frac{\Gamma}{\Gamma^2 + w^2} \right)^2; \Gamma = \frac{V}{eR}$$

# Conclusions

- Displacement: gate-voltage dependent; proceeds in steps
- Eigenmodes: modulated by the gate; steps; sensitive to the residual stress
- Insignificant effects of the mechanical degrees of freedom on the ground state energy (Coulomb diamonds) *vs* considerable influence on the current in non-equilibrium situation.