Adiabatic Transport of Cooper Pairs in Josephson Junction Circuits

J. P. Pekola, J. M. Kivioja Low Temperature Laboratory, Helsinki University of Technology, Finland J. J. Toppari University of Jyväskylä, Finland A. O. Niskanen VTT, Espoo, Finland O.Buisson, F. Balestro, J. Claudon, F. Hekking CNRS, Grenoble, France



Kylmälaboratorio Lågtemperaturlaboratoriet Low Temperature Laboratory

Contents

- 1. Adiabatic transport (phase biased circuits, environment, Landau-Zener limitation)
- 2. Experiments on a 3-junction Cooper pair pump
- 3. Flux assisted Cooper pair "sluice"
- 4. Escape measurements by a hysteretic DC-SQUID

Adiabatic transport of Cooper pairs

If the Hamiltonian of the system varies slowly enough, we may study its behaviour using the adiabatic approximation of quantum mechanics.

 $\hat{H} = \hat{H}_C + \hat{H}_J$ varies due to gate (or flux) operations.

Charge pumped in a loop of control parameters is then

$$Q_P = 2\hbar \Im m \left[\sum_{n \neq m} \oint \frac{(\hat{I}_l)_{mn}}{E_m - E_n} \langle n | \partial_{\vec{q}} m \rangle \cdot d\vec{q} \right]$$
(*m* = 0 for transport in the ground state)

(J. P., J. Toppari, M. Aunola, M. Savolainen, D. Averin, PRB **60**, 9931 (1999).)

 Q_P is related to the Berry phase γ associated with the adiabatic loop as $Q_P/(2e) = -\frac{d\gamma}{d\varphi}$ (M. Aunola and J. Toppari cond-mat/0303176)

Perfectly phase-biased adiabatic CPP





$$Q_P = 2\hbar \Im \operatorname{m} \left[\sum_{n \neq m} \oint \frac{(\hat{I}_l)_{mn}}{E_m - E_n} \langle n | \partial_{\vec{q}} m \rangle \cdot d\vec{q} \right]$$

 $Q_{\rm P}/(2e) \simeq 1 - 9E_{\rm J}/E_{\rm C}\cos{\varphi}$ (triangular path)

Adiabatic transport vs. Landau-Zener band-crossing

Adiabatic approximation is valid, if

$$|\langle 1|\dot{0}\rangle| << (E_1 - E_0)/\hbar$$

Band-crossing (LZ) becomes important, if this condition is violated.

For periodic gate operation this happens at r^2

$$f_{\rm LZ} \sim \frac{E_{\rm J}}{\hbar E_{\rm C}}$$

$H = H_{\rm J} + H_{\rm C}$



Phase fluctuations due to environment

FDT: $J(t) = 2 \int_0^\infty \frac{d\omega}{\omega} \frac{\operatorname{Re} Z_t(\omega)}{R_K}$ $\times \left\{ \operatorname{coth} \left(\frac{\hbar \omega}{2k_B T} \right) [\cos(\omega t) - 1] - i \sin(\omega t) \right\}$ $\Delta \varphi \equiv \varphi(t) - \varphi(0) \qquad \langle (\Delta \varphi)^2 \rangle = -2 \operatorname{Re} J(t)$

J. P. and J. Toppari, PRB 64, 172509 (2001) J. Martinis, S. Nam, J. Aumentado, K. Lang, C. Urbina (2003), discusses general noise-induced decoherence

Phase fluctuations due to environment



Simple idea:

Determine the fluctuations of **j** (in real time).

If $\langle (\mathbf{Dj})^2 \rangle \langle (\mathbf{p}/2)^2$, pumping is coherent

If $\langle (\mathbf{Dj})^2 \rangle \gg (\mathbf{p}/2)^2$, pumping is incoherent

In the latter case:

$$Q_{\rm P}/2e = 1 - 9E_{\rm J}/E_{\rm C}\langle\cos\varphi\rangle \to 1$$



Phase fluctuations: results in resistive environment



Pumping regimes determined by Landau-Zener bandcrossing and phase fluctuations



SQUID amplifier or escape junction for inductive read-out

Current measurement by:(a) SQUID amplifier(b) Biased Josephson junction





 $Q_{\rm P}/(2e) = 1 - [9E_{\rm J}/E_{\rm C} - (2\pi^2/3)LE_{\rm J}/\Phi_0^2]\cos\varphi$



R. Fazio, F. Hekking, J.P., preprint (2003)

Inductance limited phase fluctuations (overdamped case)



Pumping of single electrons in a 3-junction pump





H. Pothier, P. Lafarge, C. Urbina, D. Esteve, M. Devoret, EPL 17, 249 (1992).

Metrological charge pump

Accuracy of electron counting using a 7-junction electron pump

Mark W. Keller^{a)} and John M. Martinis National Institute of Standards and Technology, Boulder, Colorado 80303

Neil M. Zimmerman National Institute of Standards and Technology, Gaithersburg, Maryland 20899

Andrew H. Steinbach^{b)} National Institute of Standards and Technology, Boulder, Colorado 80303

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20

 $1/T_{mc}(1/K)$

10

0

30

Appl. Phys. Lett. 69, 1804 (1996).

The first 3 junction CPP



L. Geerligs et al., Z. Phys. B: Condensed Matter 85, 349 (1991).

R-pumps, Zorin, Lotkhov, Bogoslovsky, Niemeyer (2000, 2001)



Current / (pA)



Cooper pair pump







Experimental determination of the gate dependences



Gate modulation at different bias points



V = 21 mV



 $V = 96 \, {\rm mV}$



V = 66 mV





V = 126 mV

Pumping trajectories:



IV-curves at various frequencies



Frequency dependence of "plateaus"



Fits to LZ-crossing model yield $f_{LZ} = 20...25$ MHz, we would predict 32 MHz

Flux-assisted Cooper pair sluice



A. O. Niskanen, J. P. and H. Seppä, submitted 2003

Predicted accuracy of the (ideal) device



To study also the non-adiabaticity errors, we obtained these results by solving Schrödinger eq. and integrating in time, not by adiabatic approximation.

ω

Influence of residual Josephson coupling and offset charge



Moderate phase noise averages part of the $E_{\rm J}^{\rm res}$ -error out.



current in coil (mA)

Tunable SQUIDs?



This is more difficult to control

Apply the scheme of Romito, Plastina and Fazio?

Escape measurements using trapezoidal current pulses







Width of the escape histogram



F F

Quantronium

D. Vion, A. Aassime, A. Cottet, P. Joyez, H. Pothier, C. Urbina, D. Esteve, and M. H. Devoret, Science **296**, 886 (2002).



Possible to use an escape junction as a 'classical' ammeter.

Replacing the JJ by a DC-SQUID allows fast measurements of the state of the quantum system (O. Buisson, F. Balestro, J.P., F. Hekking, PRL, to appear.)

Pulse

measurements onDC-SQUIDs toextract phasefluctuations andcoherentcorrections to Q_P



Summary

- 1. In (conventional) Cooper pair pumps it is hard to control quasiparticle populations, supercurrent leakage, coherent corrections to quantized current and Landau-Zener transitions. Pumping cycles are complicated.
- 2. We propose and are testing a flux-assisted single island Cooper pair pump ("sluice").
- 3. Inductive measurement scheme avoids phase fluctuations: experiments on a hysteretic DC-SQUID in macroscopic quantum tunnelling regime were presented. We plan to use this technique for "noninvasive" measurements of coherent effects in Cooper pair circuits.

Adiabatic manipulation of Cooper pairs in arrays of JJs



Adiabatic quantum computing, D. Averin (1998).

Topologically stable Josephson qubits, Ioffe, Feigelman (2002).



Cooper pair manipulations in geometric qubits, Catania group (2000-).

Charging energy

Cooper pair pump $H_{c} = (2Ec/3)[(N_{1}-n_{1})^{2} + (N_{2}-n_{2})^{2} + (N_{1}-n_{1})(N_{2}-n_{2})]$ Ground state energy in $n_{v}n_{2}$ -plane





 $E_{_{\rm J}}/E_{_{\rm C}} << 1$

Coherent vs incoherent Cooper pair pumping





S. Lotkhov, S. Bogoslovsky, A. Zorin, J. Niemeyer, APL 78, 946 (2001).

Coherent pump

Circuit models of SQUID amplifier and escape measurement



Parameters of the planned measurement

Pumping frequency <100 MHz, current <30 pA (aluminium)

Niobium would allow > 1 GHz and > 0.3 nA (?)

Measurement bandwidth > $\tau_{\phi}^{-1} \sim 1 \text{ MHz}$

