

2D
holes
&
2D
electrons
in GaAs

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2D
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in Si

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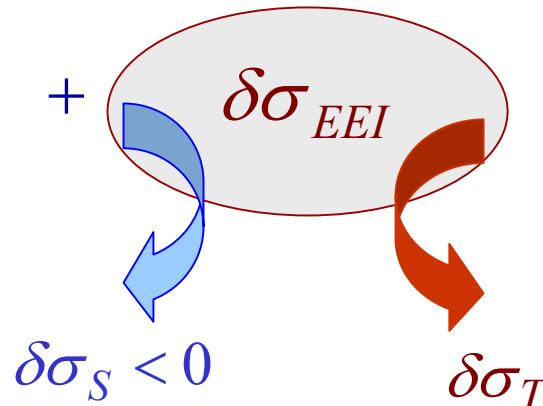
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Disorder and Interactions

Total quantum correction to the Drude conductivity:

$$\delta\sigma(T) = \delta\sigma_{WL} +$$

$$< 0$$



depends on the
interaction constant F

For large $r_s = \frac{U_C}{E_F} \propto \frac{m^*}{n^{1/2}}$ the value of F is not known.

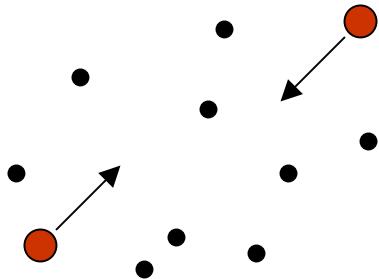
Electron-electron interactions in the diffusive regime, $k_B T \tau / \hbar < 1$

Altshuler, Aronov;
Finkelstein.

Quasi-particle interaction time, $\Delta\tau = \hbar/(k_B T)$,

is larger than the momentum relaxation time, τ

: $\Delta\tau > \tau$.



Diffusive motion enhances the interaction strength:

$$\tau_{ee}^{-1}(T) \propto T, \quad \text{instead of} \quad \tau_{ee}^{-1}(T) \propto T^2.$$

- *Logarithmic* correction to the conductance:

$$\delta\sigma(T) = \frac{e^2}{2\pi^2 \hbar} \left(1 - \frac{3}{4} F\right) \ln T \tau$$

Interactions in 2D systems in the ballistic regime

A.K.Savchenko

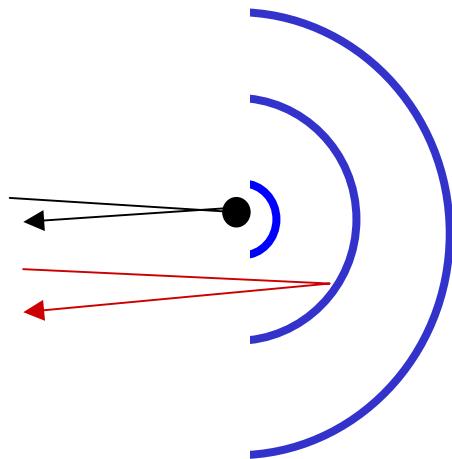
School of Physics, University of Exeter, Exeter EX4 4QL, UK

- Diffusive and ballistic regimes: $T\tau < 1$ and $T\tau > 1$.
- { • Hole-hole interactions in a 2DHG in *GaAs* and ‘metallic’ $\rho_{xx}(T)$ with $d\rho_{xx}/dT > 0$.
- { • Electron interactions in a 2DEG in *Si* in the ballistic regime.
- { • Electron interactions in a 2DEG in *GaAs* in the ballistic regime in the presence of a long-range fluctuation potential.

Interaction corrections to σ_D in the ballistic regime $T\tau > 1$.

Zala, Narozhny, Aleiner (2001)

Physics: Coherent electron scattering by *Friedel oscillation* produced by a *point scatterer*.



$$\delta\rho(r) \propto \frac{1}{r^2} \exp(i2k_F r)$$

Phase difference is cancelled by the phase of the Friedel oscillation \Rightarrow interference.

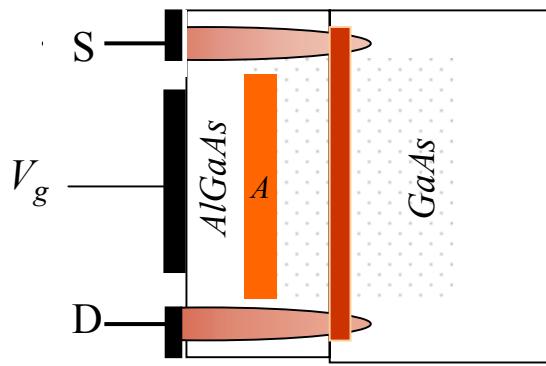
- Prediction:

$$\delta\sigma(T) = \frac{e^2}{\pi\hbar} \left(1 + \frac{3F_o^\sigma}{1 + F_o^\sigma} \right) \cdot T\tau$$

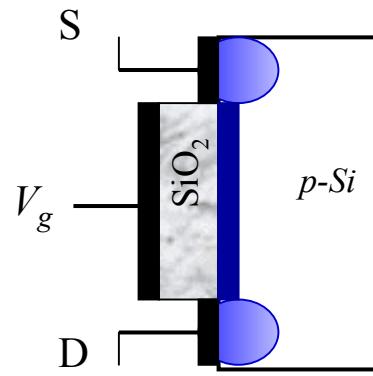
- linear correction to the conductance.

2D structures

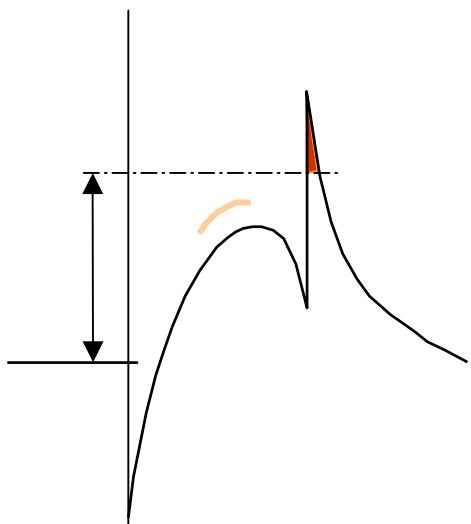
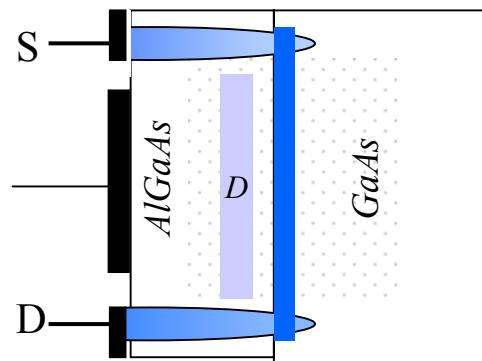
2DHG in
GaAs/AlGaAs



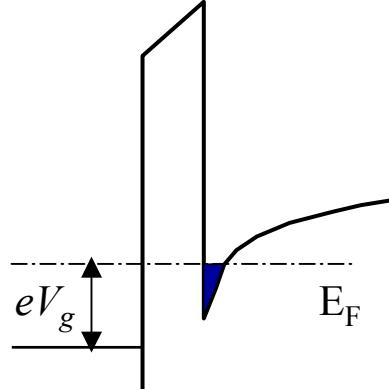
2DEG in Si
MOSFET



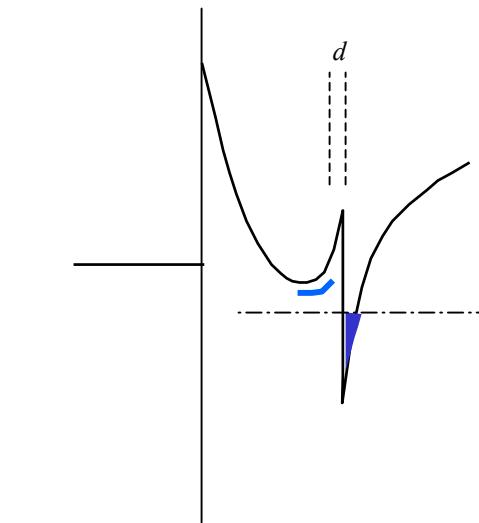
2DEG in
GaAs/AlGaAs



Point-like scatterers
(background impurities)

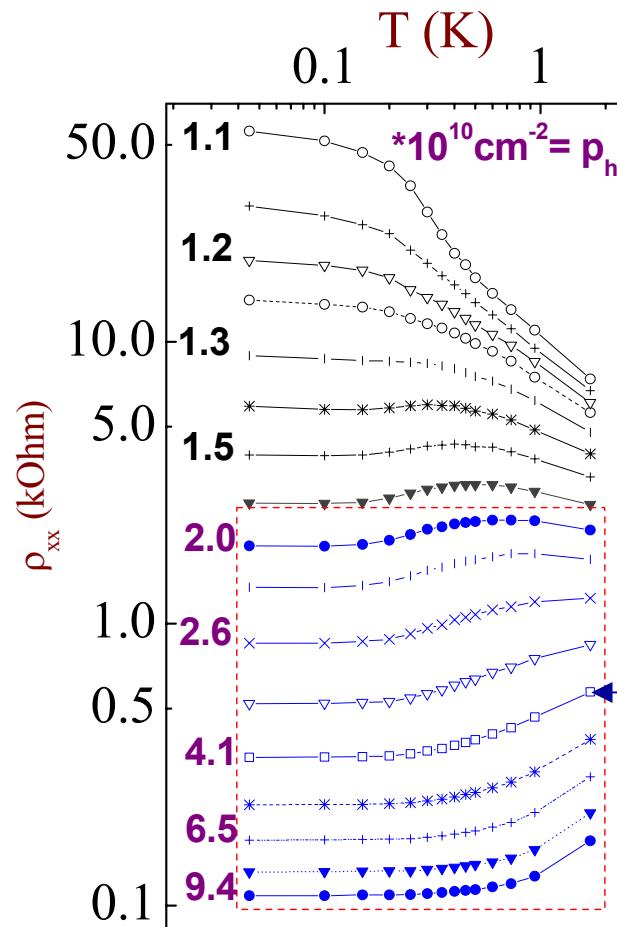


Point-like scatterers



Long-range potential
(correlation length d - spacer thickness)

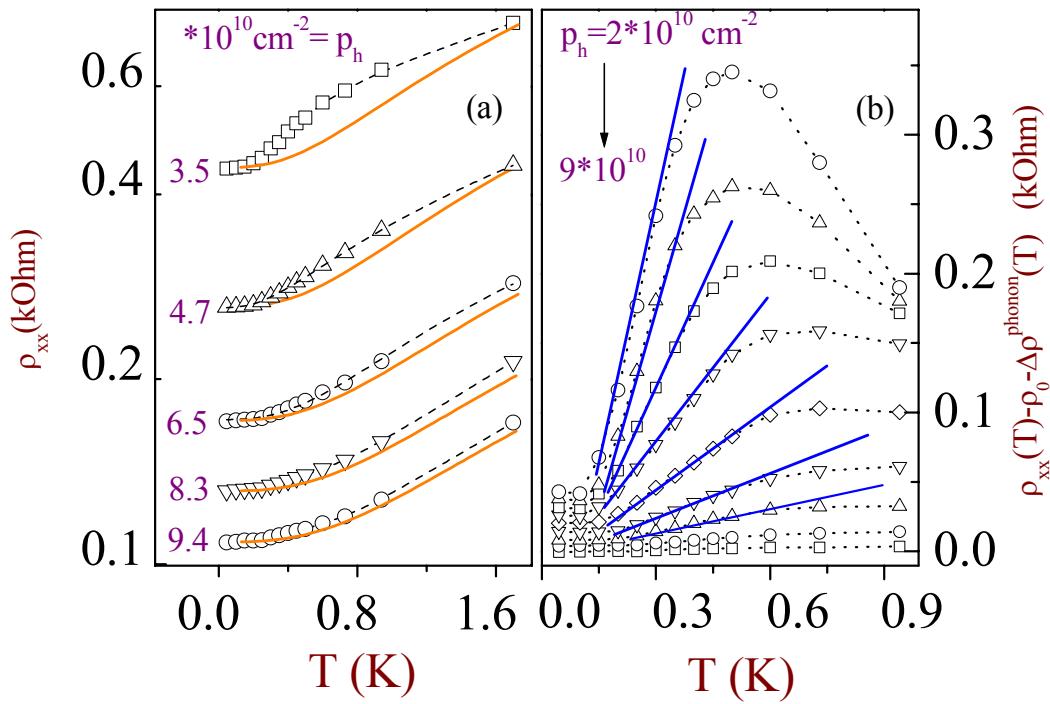
Temperature dependence of resistivity near the crossover.



2DHG in GaAs/AlGaAs

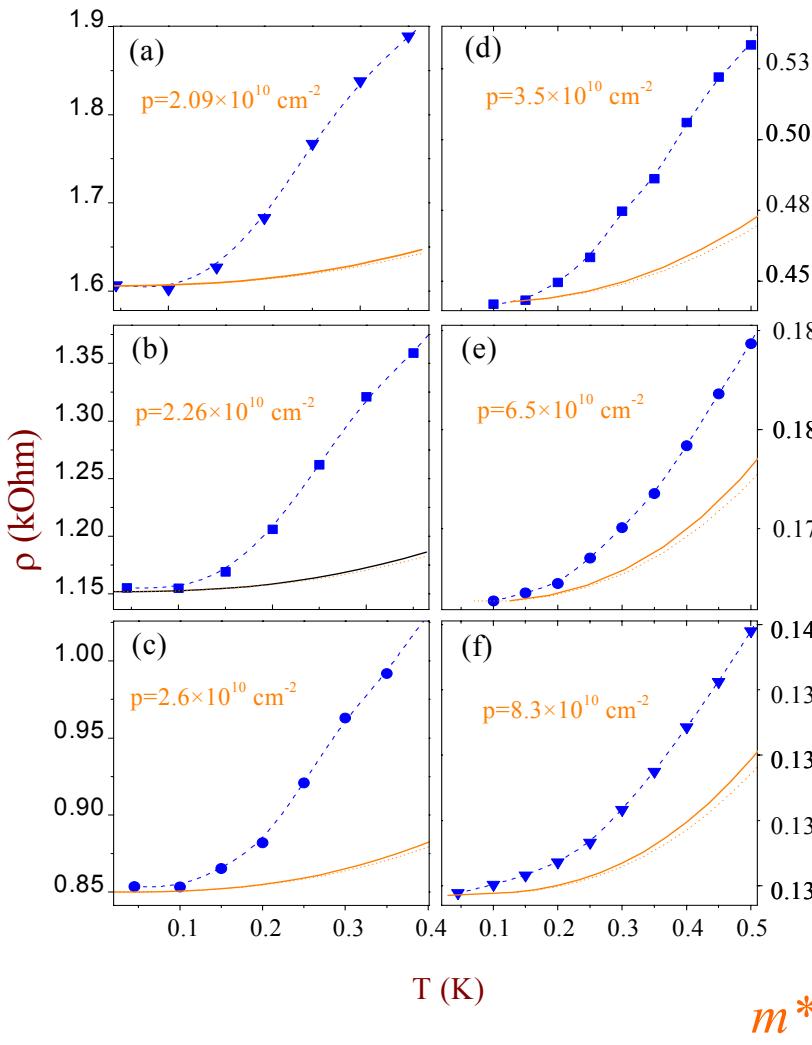
‘metallic’ region

Phonon and impurity scattering in ‘metallic’ $R(T)$

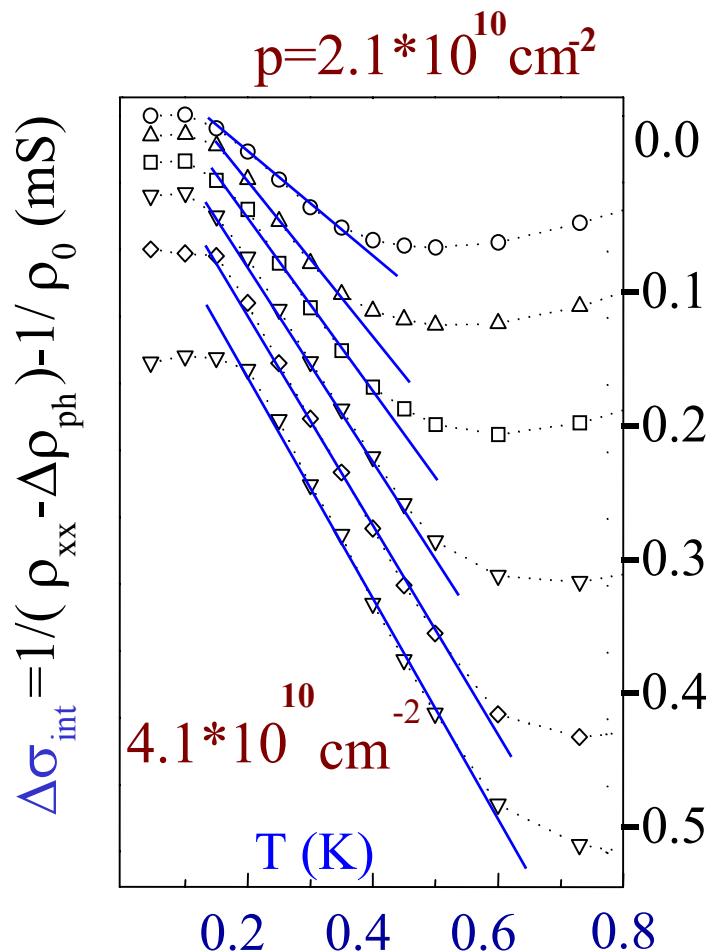


*Theory of phonon scattering:
Karpus (1990)*

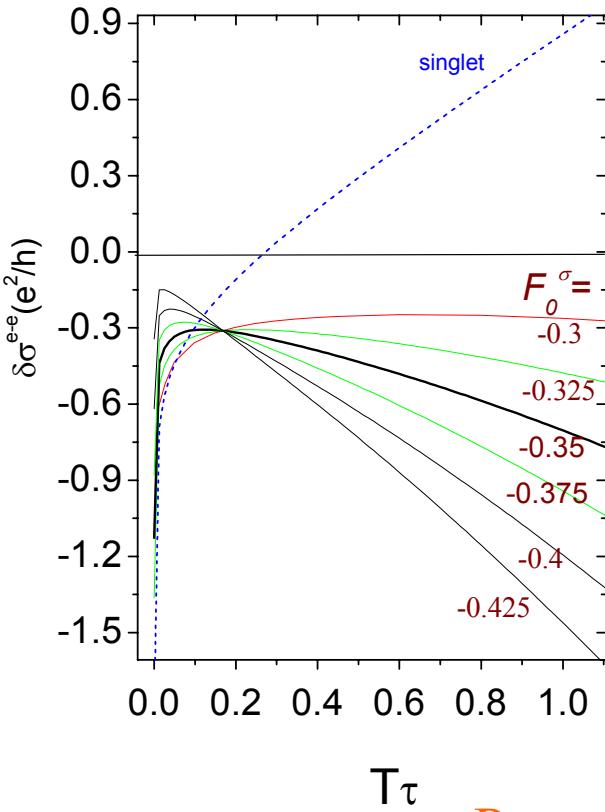
The phonon contribution to $\rho_{xx}(T)$.



Interaction correction in the *ballistic* regime, $k_B T \tau / \hbar > 1$.



Zala, Narozhny, Aleiner (2001)



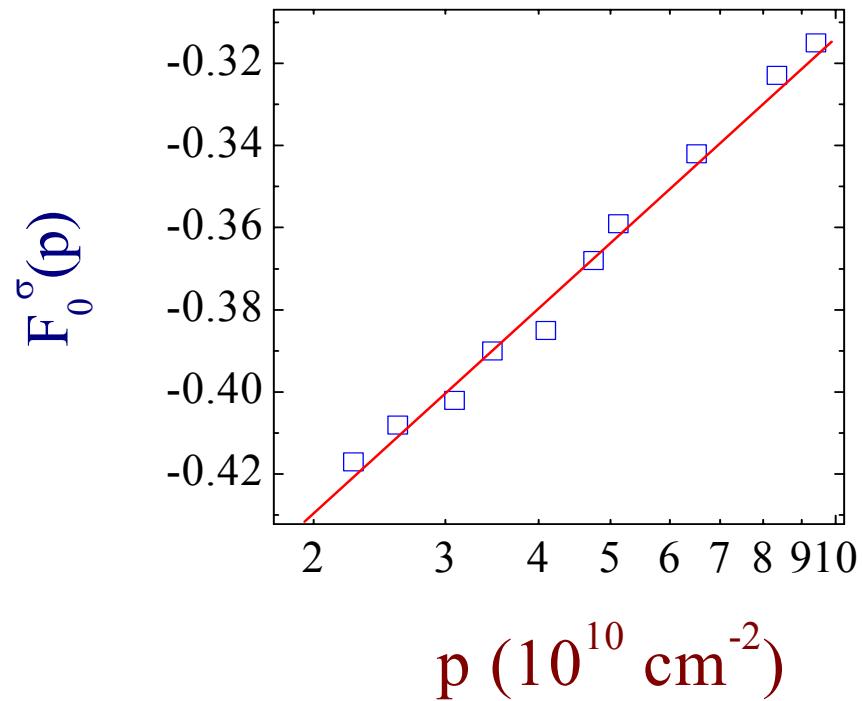
$$\delta\sigma(T) = \frac{e^2}{\pi\hbar} \frac{k_B T \tau}{\hbar} \left(1 + \frac{3F_o^\sigma}{1 + F_o^\sigma} \right)$$

B

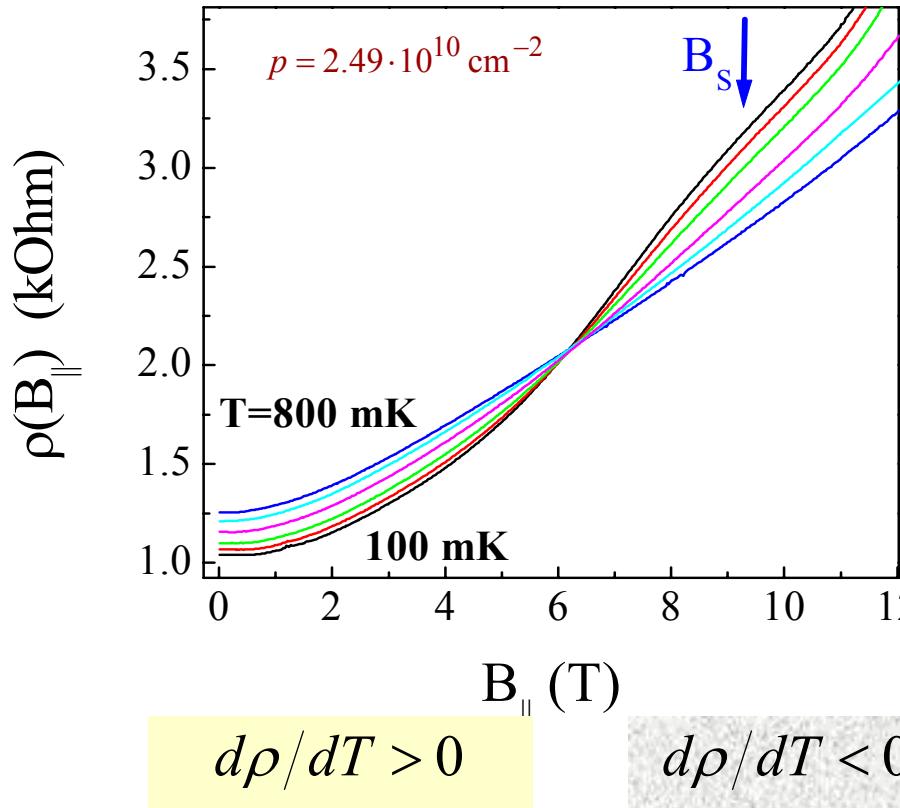
$$= \sigma_o \left(1 + \cancel{\frac{3F_o^\sigma}{1 + F_o^\sigma}} \right) \cdot \frac{k_B T}{E_F}$$

singlet *triplet*

Interaction constant F as a function of the hole density



Positive magnetoresistance in parallel field



B_s is the field of full spin polarisation, $B_s = 2E_F/g^* \mu_B$.

Change in $\delta\sigma(T)$ with B_{par}

$B_{\text{par}}=0$

$$\delta\sigma(T) = \sigma_o \left(1 + \frac{3F_o^\sigma}{1 + F_o^\sigma} \right) \cdot \frac{k_B T}{E_F}$$

< 0

‘Metal’

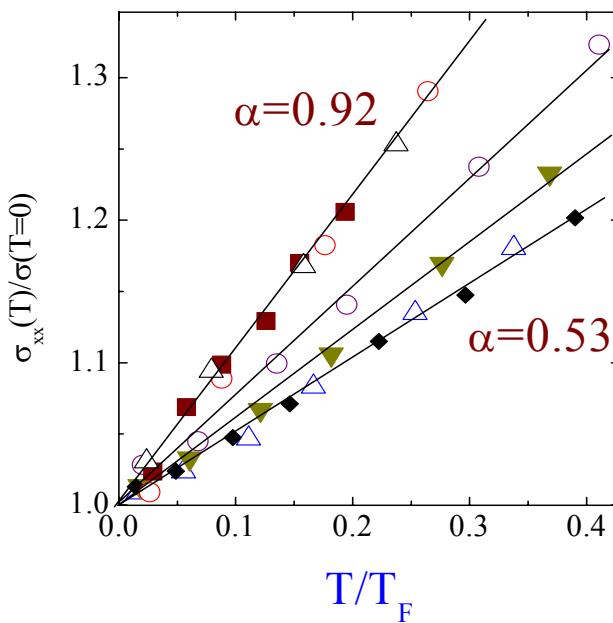
$B_{\text{par}} > B_s$

$$\delta\sigma(T) = \sigma_o \cdot \frac{k_B T}{E_F}$$

‘Insulator’

$B=B_s$

$$\delta\sigma(T) = \alpha \cdot \sigma_o \cdot \frac{k_B T}{E_F}$$

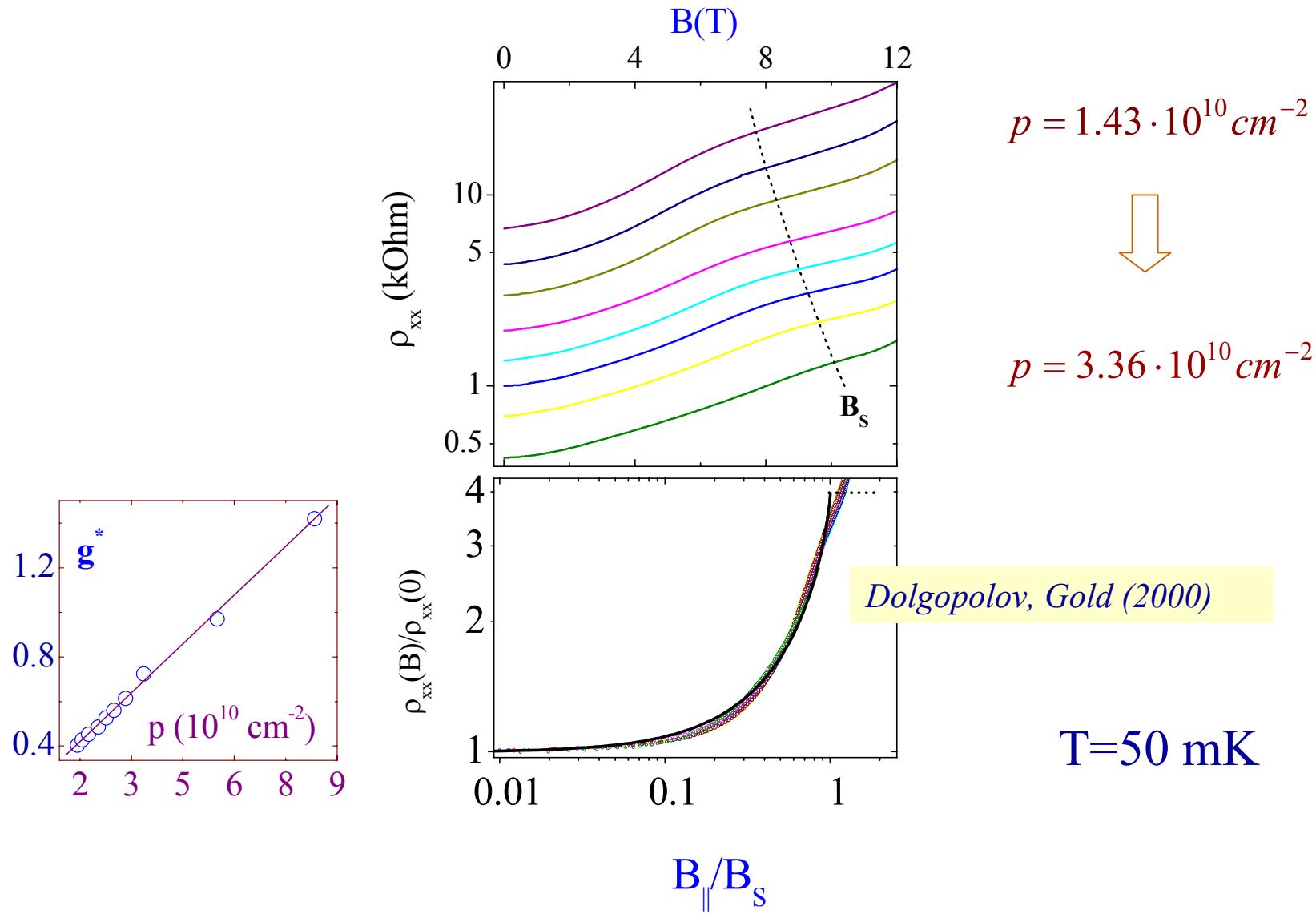


$$p = 1.43 - 1.75 \cdot 10^{10} \text{ cm}^{-2}$$

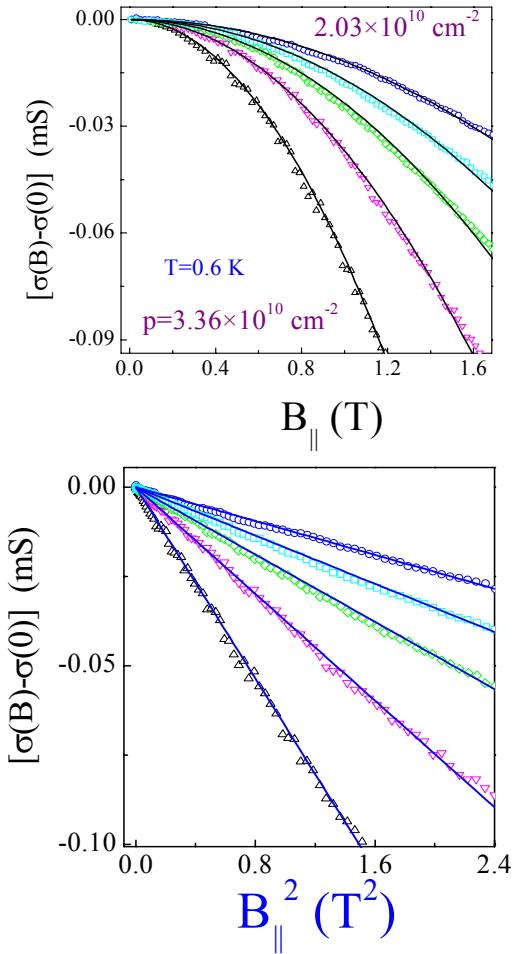


$$p = 2.49 \cdot 10^{10} \text{ cm}^{-2}$$

General view on the PMR and the g-factor



Positive magnetoresistance in small parallel field

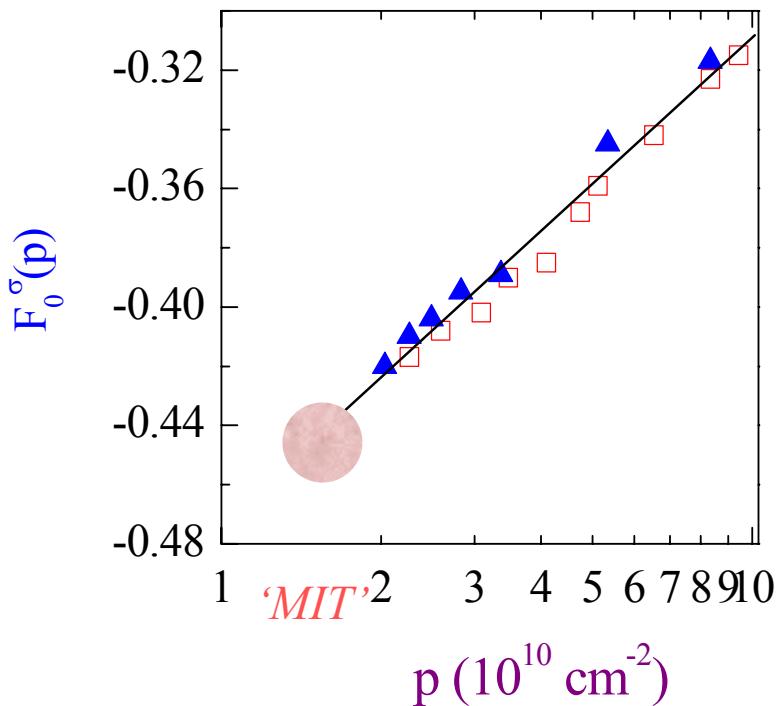


Zala, Narozhny, Aleiner (2001)

$$\text{at } \frac{g^* \mu_B B}{2T} < 1 + F_o^\sigma$$

$$\sigma(B) - \sigma(0) \approx \frac{e^2}{\pi \hbar} \frac{2F_o^\sigma}{(1 + F_o^\sigma)} \frac{T \tau}{\hbar} \frac{\left(g^* \mu_B B / 2T\right)^2}{3} f(F_o^\sigma)$$

Interaction constant in the triplet channel $F_o^\sigma(p)$.



□ *from $\sigma(T)$
at $B = 0$*

▲ *from $\sigma(B,T)$*

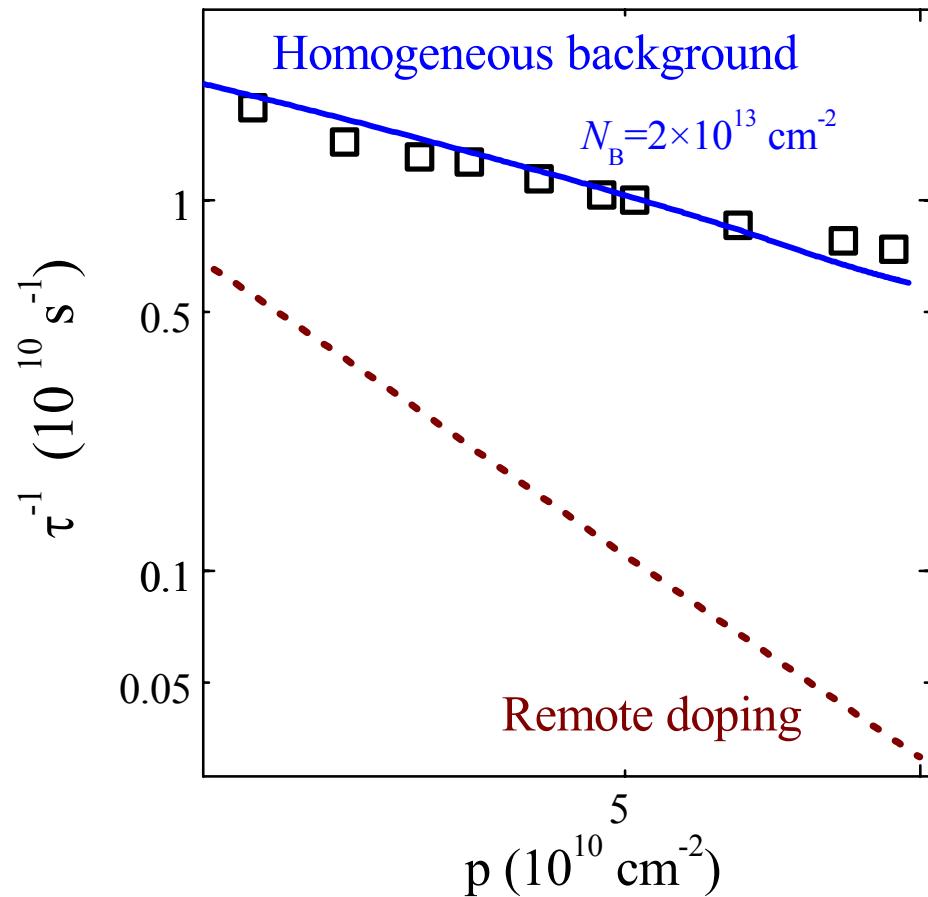
$$at \frac{g^* \mu B}{2T} < 1 + F_o^\sigma$$

Magnetic susceptibility: $\chi = \frac{\chi_o}{1 + F_o^\sigma}$

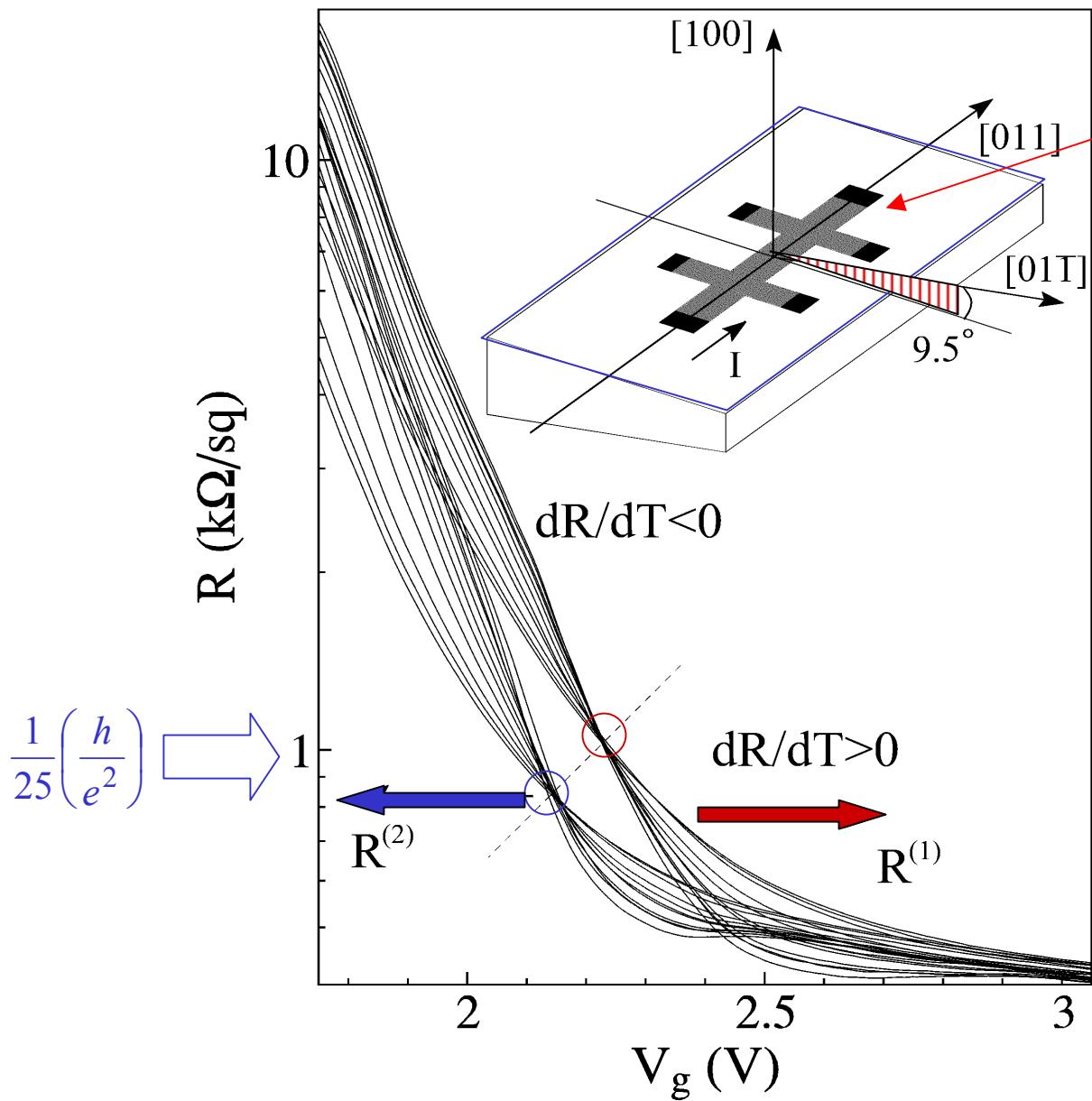
(Stoner instability $F_o^\sigma = -1$)

Proskuryakov, Savchenko, Safonov, Pepper, Simmons,
D. A. Ritchie, Phys. Rev. Lett. 89, 076406 (2002)

Short-range scattering in the 2DHG in GaAs



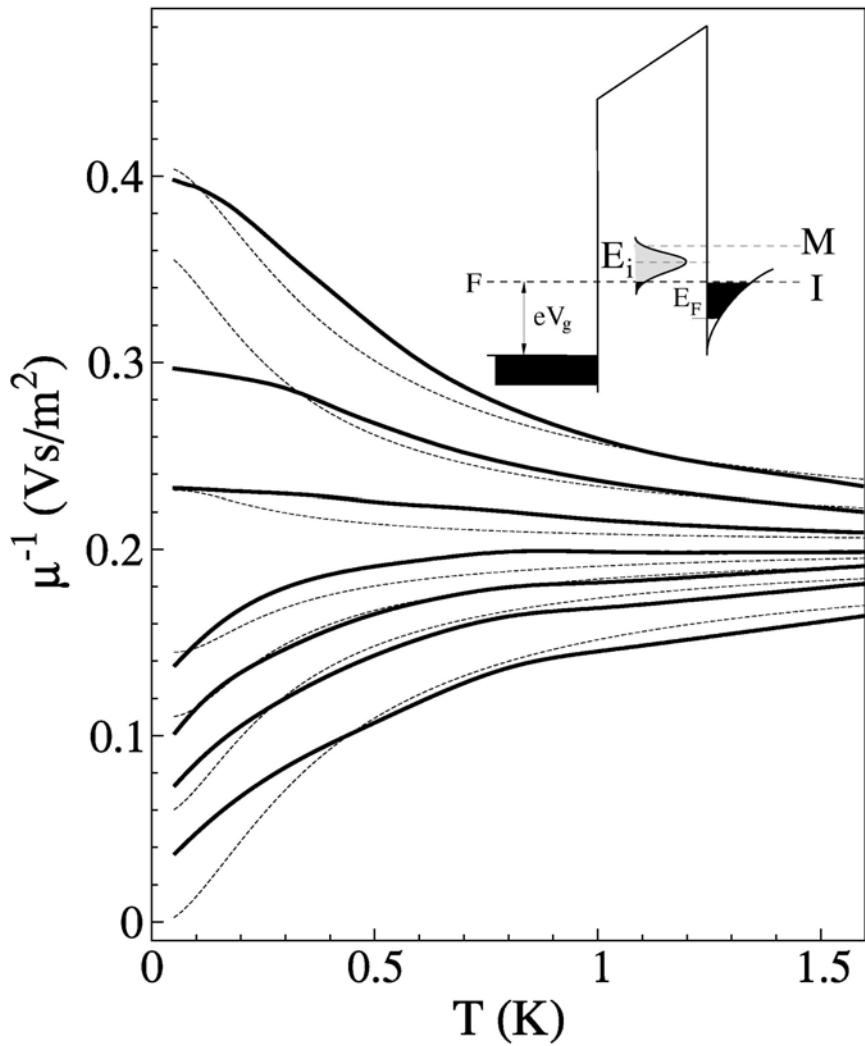
Crossover in the sign of dR/dT for different directions of V_g sweep



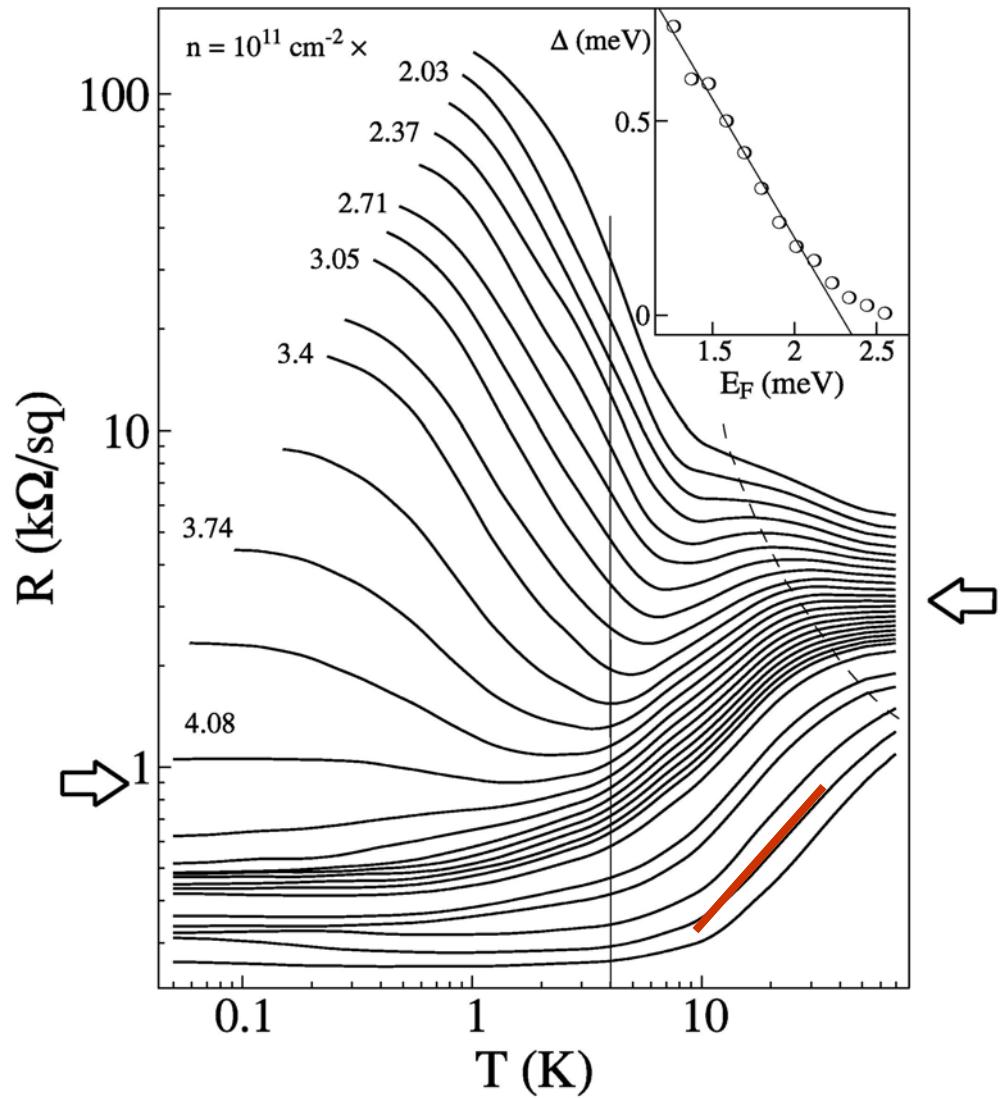
2DEG on vicinal
Si surface

Safonov, Roshko, Savchenko,
Pogosov, Kvon, PRL (2001)

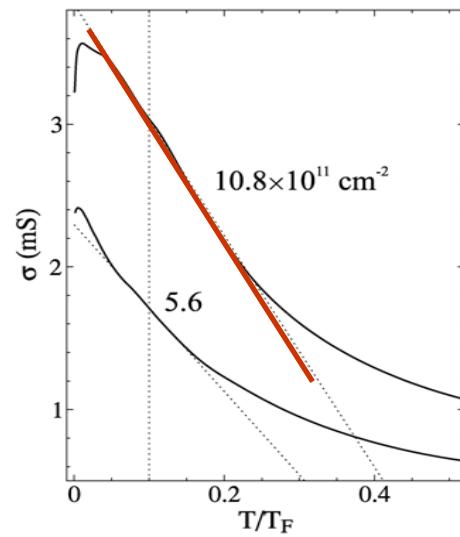
A narrow impurity band in the origin of the ‘MIT’ in the 2DEG on vicinal Si.



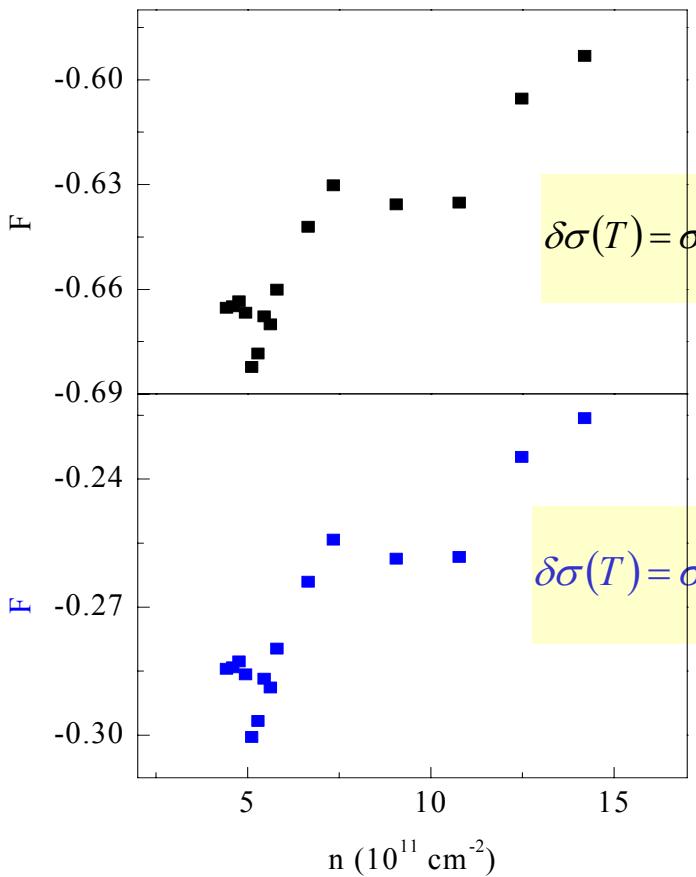
Temperature dependence of resistance of a 2DEG in vicinal Si



2DEG on vicinal
Si surface



Interaction constant $F_0^\sigma(n)$ of the 2DEG in vicinal Si

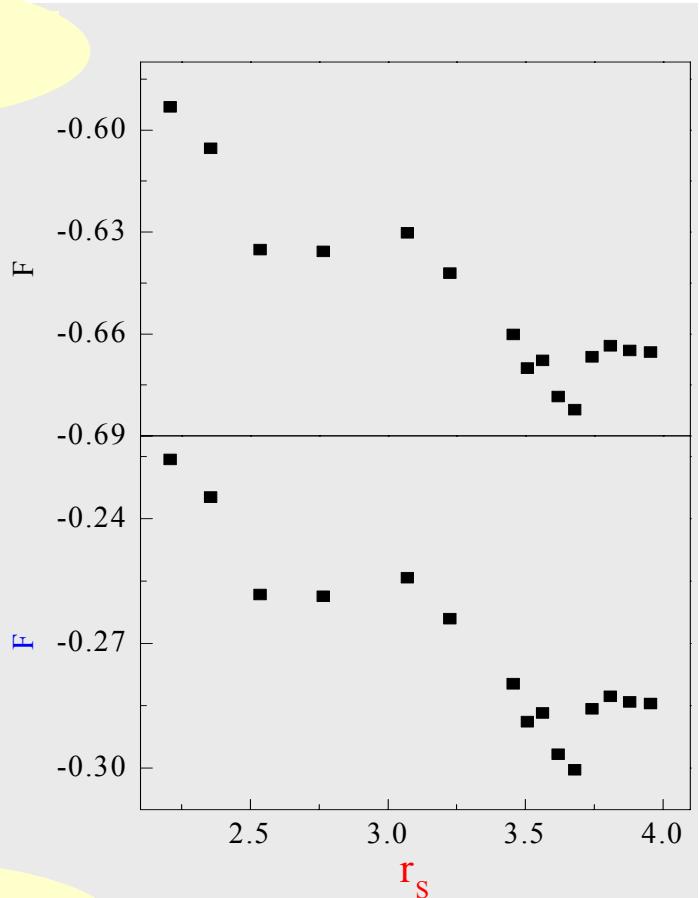


One valley

$$\delta\sigma(T) = \sigma_o \left(1 + \frac{3F_o^\sigma}{1 + F_o^\sigma} \right) \cdot \frac{k_B T}{2E_F}$$

$$\delta\sigma(T) = \sigma_o \left(1 + \frac{15F_o^\sigma}{1 + F_o^\sigma} \right) \cdot \frac{k_B T}{2E_F}$$

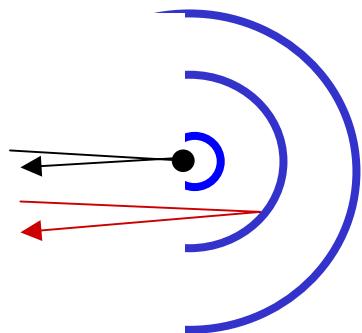
Two valleys



Ballistic regime, $T\tau > 1$, in short- and long-range fluctuation potential

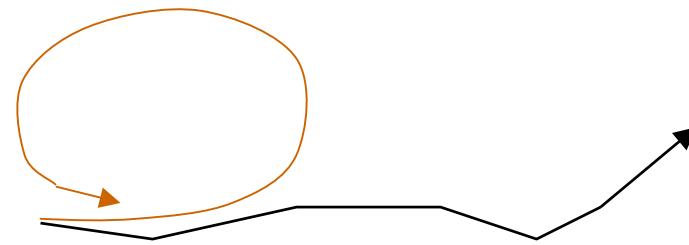
Short-range potential

(essential for *Zala, Narozhny, Aleiner (2001)*)



Long-range potential

(discussed by *Gornyi and Mirlin (2002)*)



- Smooth scattering potential **suppresses** interaction correction (at $B=0$) .

Measurements of the **interaction correction** by the **parabolic NMR** was discussed in relation to the *diffusive* regime:

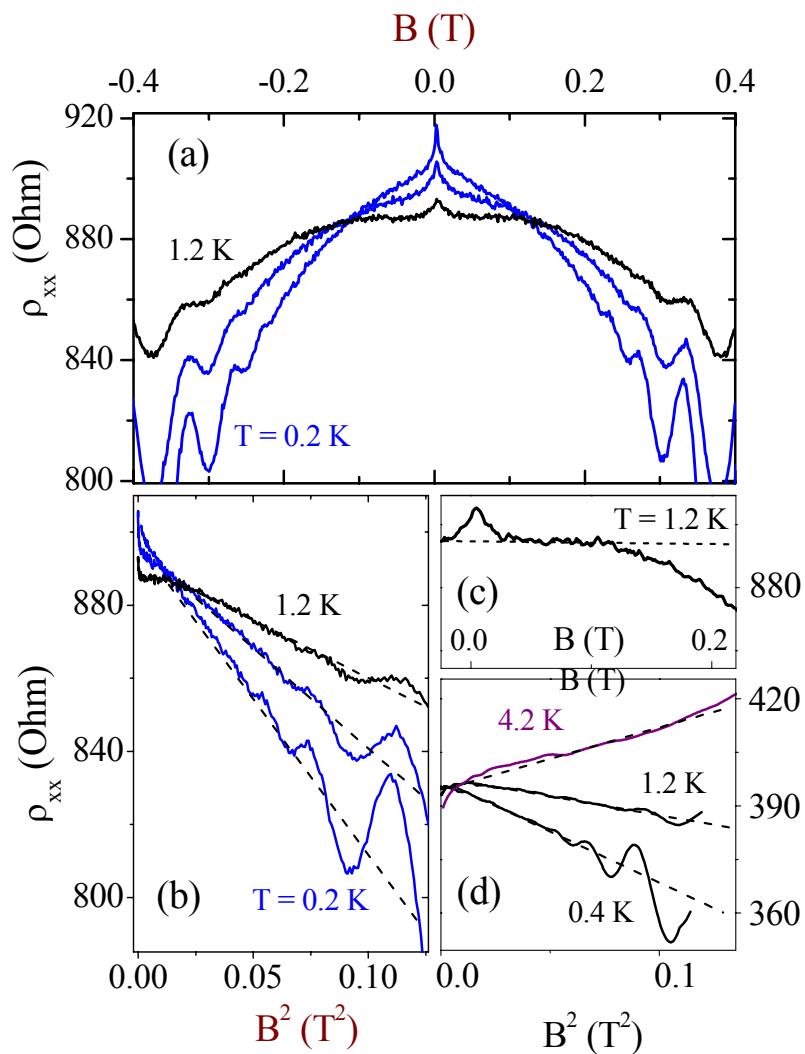
Paalanen, Tsui, Hwang, PRL (1983).

- Strong field, $\omega_c\tau > 1$, **restores** interactions:

$$\rho_{xx} = \frac{1}{\sigma_0} + \frac{1}{\sigma_0^2} (\omega_c\tau)^2 \delta\sigma_{xx}^{ee}(T)$$

Electron interactions in the ballistic regime *in the long-range fluctuation potential*

Theory: *Gornyi, Mirlin, PRL (2003)*



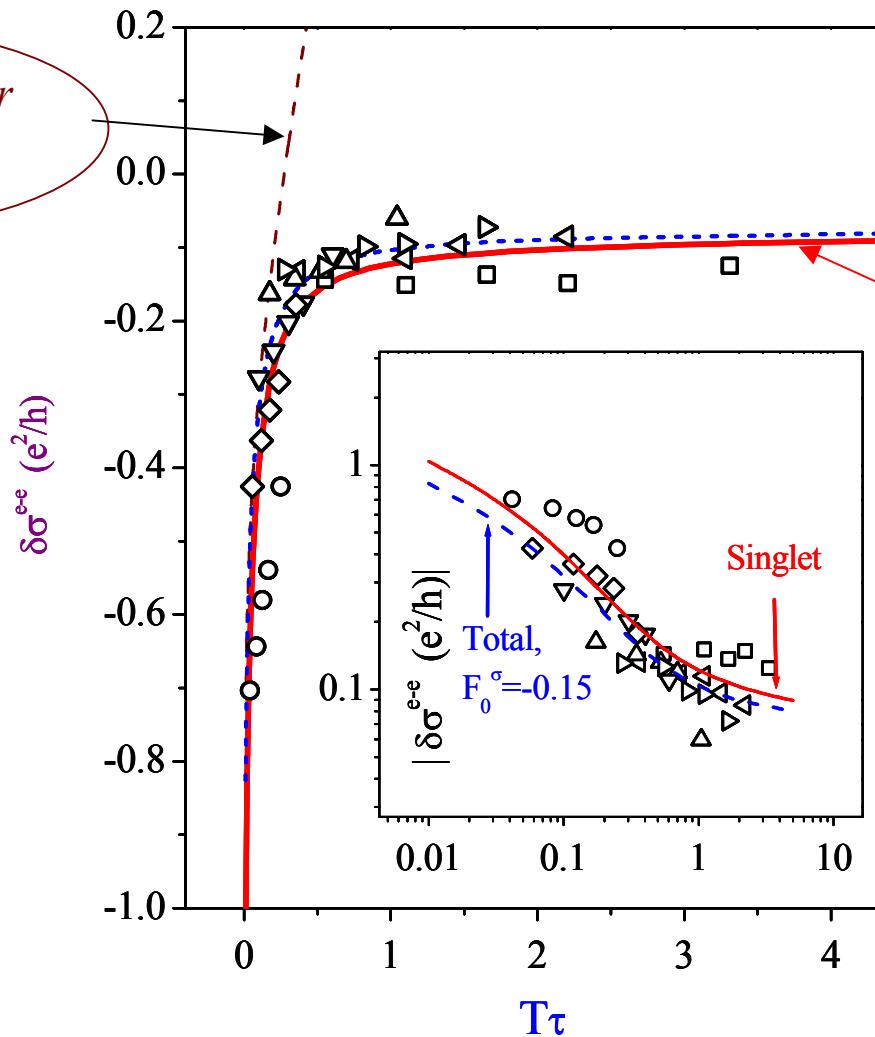
Experiment:

Li, Proskuryakov, Savchenko, Linfield, Ritchie, PRL (2003)

2DEG in GaAs/AlGaAs

Interaction correction $\delta\sigma^{ee}(T)$ in strong magnetic fields

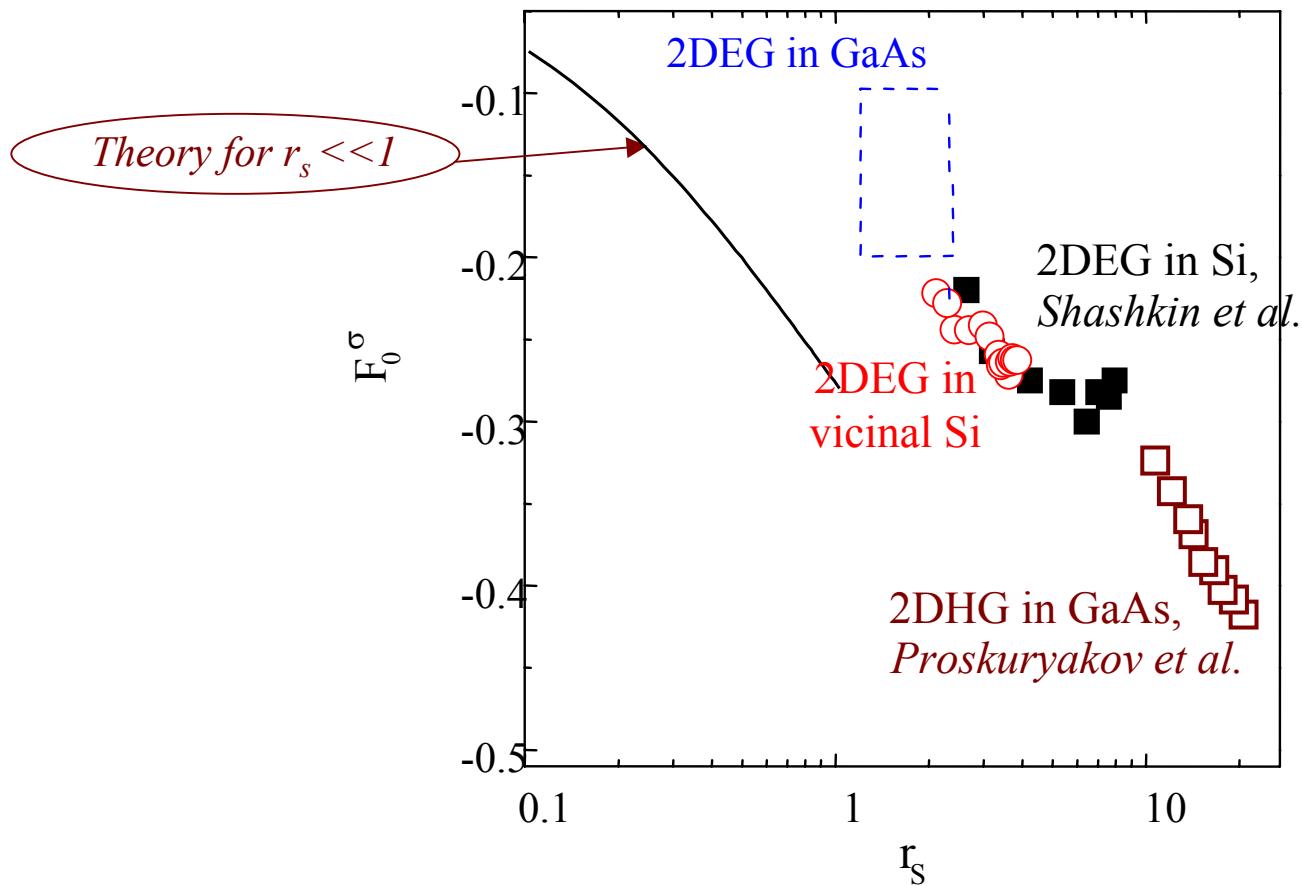
Zala, Narozhny, Aleiner
(singlet) : point-like scatterers, $B=0$



Gornyi, Mirlin (singlet):
long-range potential

$$\rho_{xx} = \frac{1}{\sigma_0} + \frac{1}{\sigma_0^2} \left(\mu^2 B^2 \right) \delta\sigma_{xx}^{ee}(T)$$

Interaction parameter F_o^σ as a function of r_s in different 2D systems



Interaction parameter F_0^σ as a function of r_s in different 2D systems

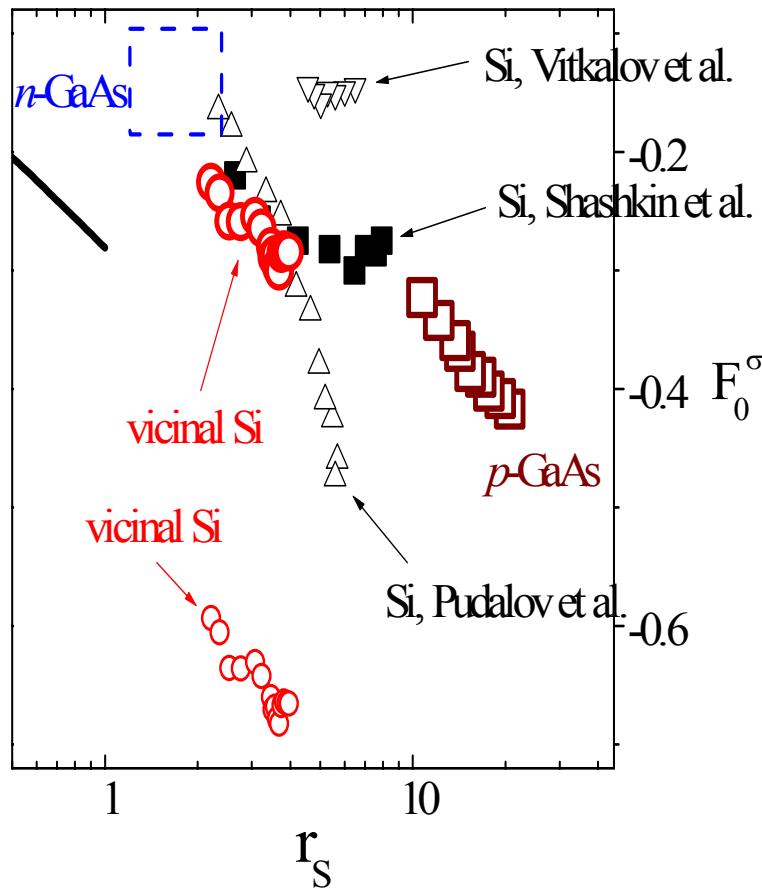


Fig 1

Conclusions

- Interaction effects in the *ballistic regime* contribute to the ‘metallic’ $\rho_{xx}(T)$ in the systems with short-range scatterers:
2DHG in *GaAs/AlGaAs* and 2DEG in *Si*.
- Interaction effects in the *ballistic regime* depend on the character of the fluctuation potential. They have different manifestation in the situation of long-range potential:
as a negative magnetoresistance of a 2DEG in GaAs/AlGaAs.